Approximate Analysis to the KdV-Burgers Equation

Zhaosheng Feng

Department of Mathematics University of Texas-Pan American 1201 W. University Dr. Edinburg, Texas 78539, USA E-mail: zsfeng@utpa.edu

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- KdV-Burgers Equation
- Planar Polynomial Systems and Abel Equation

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- Generalized Abel Equation
- Property of Our Operator
- Two Theorems

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- Approximate Solution to 2D KdV-Burgers Equation

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• Generalized Korteweg-de Vries-Burgers equation [1, 2]

$$u_t + \left(\delta u_{xx} + \frac{\beta}{p}u^p\right)_x + \alpha u_x - \mu u_{xx} = 0, \tag{1}$$

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where *u* is a function of *x* and *t*, α , β and p > 0 are real constants, μ and δ are coefficients of dissipation and dispersion, respectively.

• The type of such problems arises in modeling waves generated by a wavemaker in a channel and waves incoming from deep water into nearshore zones.

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- The type of such problems arises in modeling waves generated by a wavemaker in a channel and waves incoming from deep water into nearshore zones.
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$$u(x, t) = \frac{3\beta^2}{25\alpha s} \operatorname{sech}^2 \Psi - \frac{6\beta^2}{25\alpha s} \tanh \Psi \pm \frac{6\beta^2}{25\alpha s}, \qquad (5)$$

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where

$$\Psi = \left[\frac{1}{2}\left(-\frac{\beta}{5s}x \pm \frac{6\beta^3}{125s^2}t\right)\right].$$

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[6] R.S. Johnson, J. Fluid Mech. 42 (1970), 49-60.

[7] Z. Feng, J. Phys. A (Math. Gen.) 36 (2003), 8817-8827.

[8] Z. Feng, Nonlinearity, 20 (2007), 343–356.

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 Choices of α = 0 and p = 2 lead equation (1) to the standard form of the Korteweg-de Vries-Burgers equation [6]:

$$u_t + \alpha u u_x + \beta u_{xx} + s u_{xxx} = 0. \tag{4}$$

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• Solitary wave solutions of equation (4) are as follows [7, 8, 9]:

$$u(x, t) = \frac{3\beta^2}{25\alpha s} \operatorname{sech}^2 \Psi - \frac{6\beta^2}{25\alpha s} \tanh \Psi \pm \frac{6\beta^2}{25\alpha s}, \qquad (5)$$

where

$$\Psi = \left[\frac{1}{2}\left(-\frac{\beta}{5s}x \pm \frac{6\beta^3}{125s^2}t\right)\right]$$

- [6] R.S. Johnson, J. Fluid Mech. 42 (1970), 49-60.
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- [8] Z. Feng, Nonlinearity, 20 (2007), 343–356.
- [9] Z. Feng and S. Zheng, Z. angew. Math. Phys. 60 (2009), 756–773.

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Figures of Wa	ve Solutions			



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Planar Poly	nomial Systems an	nd Abel Equation		

• Consider planar polynomial systems of the form

$$\dot{x} = -y + p(x, y), \quad \dot{y} = x + q(x, y)$$
 (6)

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Planar Poly	vnomial Systems ar	nd Abel Equation		

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with homogeneous polynomials p(x, y) and q(x, y) of degree k.

• For the Poincaré center problem, setting $x = r \cos \theta$, $y = r \sin \theta$ gives

$$\frac{dr}{d\theta} = \frac{r^k \xi(\theta)}{1 + r^{k-1} \eta(\theta)},\tag{7}$$

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where ξ and η are polynomials in $\cos \theta$ and $\sin \theta$ of degree k + 1.

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Planar Polynomial Systems and Abel Equation

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Planar Polynomial Systems and Abel Equation

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• Making the coordinate transformation

$$\rho = \frac{r^{k-1}}{1 + r^{k-1}\eta(\theta)}$$

we get an Abel equation

$$\frac{d\rho}{d\theta} = a(\theta)\rho^2 + b(\theta)\rho^3,$$

where $a = (k-1)\mathcal{E} + n'$ and $b = (1-k)\mathcal{E}n$.

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Traveling Way	ve Solution			

• Assume that equation (1) has the traveling wave solution of the form

$$u(x,t) = u(\xi), \quad \xi = x - ct,$$

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Introduction	Qualitative Analysis	Approximate Solution	Conclusion O	Acknowledgement O
Traveling Way	ve Solution			

• Assume that equation (1) has the traveling wave solution of the form

$$u(x,t) = u(\xi), \quad \xi = x - ct,$$

where $c \neq 0$ is the wave velocity. Then equation (1) becomes

$$\delta u''' - \mu u'' + (\alpha - c)u' + \beta u^{p-1}u' = 0, \tag{8}$$

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where $u' = du/d\xi$. Integrating equation (8) once gives

$$u'' - gu' - eu - fu^p - d = 0,$$
(9)

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| Traveling W | Vave Solution | | | |

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where $e = \frac{c-\alpha}{\delta}$, $g = \frac{\mu}{\delta}$, $f = -\frac{\beta}{p\delta}$ and *d* is an integration constant.

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Traveling W	Vave Solution			

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$$u'' - gu' - eu - fu^p - d = 0, (9)$$

where $e = \frac{c-\alpha}{\delta}$, $g = \frac{\mu}{\delta}$, $f = -\frac{\beta}{p\delta}$ and *d* is an integration constant.

• Assume that y = u and u' = z, then equation (9) is equivalent to

$$\begin{cases} y' = z, \\ z' = ey + gz + fy^p + d. \end{cases}$$
(10)

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Global Struct	ure of $p = 2$			



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Transformed to Abel Equation						

$$\frac{dz}{dy} = \frac{ey + gz + fy^p + d}{z}.$$
(11)

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Introduction	Qualitative Analysis	Approximate Solution	Conclusion O	Acknowledgement O
Transformed to Abel Equation		l		

$$\frac{dz}{dy} = \frac{ey + gz + fy^p + d}{z}.$$
(11)

• Let $z = r^{-1}$. Equation (11) reduces to

$$\frac{dr}{dy} = a(y)r^2 + b(y)r^3,$$
(12)

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Introduction	Qualitative Analysis	Approximate Solution	Conclusion O	Acknowledgement O
Transforme	ed to Abel Equation			

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where a(y) = -g and $b(y) = -(ey + fy^p + d)$.

Introduction 0000000	Qualitative Analysis	Approximate Solution	Conclusion O	Acknowledgement O
Transformed	l to Abel Equation			

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where a(y) = -g and $b(y) = -(ey + fy^p + d)$.

• Question: Under what condition one can determine the number of closed solutions of the Abel equation (12).

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Transformed to Abel Equation				

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• Let $z = r^{-1}$. Equation (11) reduces to

$$\frac{dr}{dy} = a(y)r^2 + b(y)r^3, \tag{12}$$

where a(y) = -g and $b(y) = -(ey + fy^p + d)$.

- Question: Under what condition one can determine the number of closed solutions of the Abel equation (12).
- Open Problem: There have been two longstanding problems, called the Poincaré center-focus problem and the local Hilbert 16th problem. Both are closely related to the Bautin quantities and the Bautin ideal of the Abel equation.

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Integral Form				

$$r' = a(t)r^2 + b(t)r^n, \quad r(t_0) = c, \quad t \in [t_0, t_1], \quad n \ge 3.$$
 (13)

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Integral Form				

$$r' = a(t)r^2 + b(t)r^n$$
, $r(t_0) = c$, $t \in [t_0, t_1]$, $n \ge 3$. (13)

• Dividing both sides of equation (13) by r^2 gives

$$\frac{r'}{r^2} = a(t) + b(t)r^{n-2}.$$
(14)

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Integral Form				

$$r' = a(t)r^2 + b(t)r^n, \quad r(t_0) = c, \quad t \in [t_0, t_1], \quad n \ge 3.$$
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$$\frac{r'}{r^2} = a(t) + b(t)r^{n-2}.$$
(14)

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• Integrating equation (14) from t_0 to t yields

$$r(t) = \frac{c}{1 - cA(t) - c \int_{t_0}^t b(\tau) r^{n-2} d\tau},$$
(15)

where $A(t) = \int_{t_0}^t a(\tau) d\tau$.

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Integral Form				

$$r' = a(t)r^2 + b(t)r^n, \quad r(t_0) = c, \quad t \in [t_0, t_1], \quad n \ge 3.$$
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(15)

where $A(t) = \int_{t_0}^t a(\tau) d\tau$. • Rewrite equation (15) as

$$r(t) = c \left(1 + A(t) + r(t) \int_{t_0}^t b(\tau) r^{n-2} d\tau \right).$$
(16)

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A Nonlinear (Operator			

 $T_c: \mathcal{C}[0,1] \to \mathcal{C}[0,1],$

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A Nonlinear Opera	ntor			

 $T_c: \mathcal{C}[0,1] \to \mathcal{C}[0,1],$

$$T_c(f)(t) \stackrel{def}{=} \frac{c}{1 - cA(t) - c \int_0^t b(\tau) f(\tau)^{n-2} d\tau},$$

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A Nonlinea	ar Operator			

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$$T_c(f)(t) \stackrel{def}{=} \frac{c}{1 - cA(t) - c \int_0^t b(\tau) f(\tau)^{n-2} d\tau},$$

for given $a, b \in C[0, 1]$ and $c \in \mathbb{R}$. Obviously, T_c is well defined on an arbitrary bounded set of C[0, 1] if c is suitably small. Let us first observe some useful properties of T_c .

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A Nonlinea	r Operator			

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Property of (Our Operator			

Lemma (1)

For $f \in C[0,1]$ and $c \in \mathbb{R}$ with $||f|| \leq M$ and $|c| < c_0 \stackrel{def}{=} (||a|| + ||b|| M^{n-2})^{-1}$, $T_c(f)$ is well defined and differentiable, and satisfies

$$\frac{d}{dt}T_c(f)(t) = a(t)[T_c(f)(t)]^2 + b(t)[T_c(f)(t)]^2 f(t)^{n-2}$$

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Property of	f Our Operator			

Lemma (1)

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$$\frac{d}{dt}T_c(f)(t) = a(t)[T_c(f)(t)]^2 + b(t)[T_c(f)(t)]^2 f(t)^{n-2}$$

Furthermore, we have an identity

$$T_c(f)(t) - T_c(g)(t) = T_c(f)(t)T_c(g)(t)\int_0^t b(\tau)(f(\tau)^{n-2} - g(\tau)^{n-2})d\tau,$$

$$0 \le t \le 1$$

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Introduction 00000000	Qualitative Analysis	Approximate Solution	Conclusion O	Acknowledgement O
Property of	f Our Operator			

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$$0 \le t \le 1$$

for arbitrary $f, g \in C[0, 1]$ and $c \in \mathcal{R}$ with $||f||, ||g|| \leq M$ and $|c| < c_0$.

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Outline of the	Proof			

Step 1: well-defined

$$1 - cA(t) - c \int_0^t b(\tau) f(\tau)^{n-2} d\tau = 0 \Rightarrow$$

$$|c| \ge \frac{1}{|A(t)| + \int_0^t |b(\tau)f(\tau)^{n-2}| d\tau} \ge \frac{1}{||a|| + ||b|| M^{n-2}}.$$

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Outline of the	Proof			

Step 1: well-defined

$$\begin{aligned} 1 - cA(t) - c \int_0^t b(\tau) f(\tau)^{n-2} d\tau &= 0 \Rightarrow \\ |c| &\ge \frac{1}{|A(t)| + \int_0^t |b(\tau) f(\tau)^{n-2}| \, d\tau} \ge \frac{1}{||a|| + ||b|| M^{n-2}}. \end{aligned}$$

Step 2: A direct calculation gives

$$\frac{d}{dt}T_{c}(f)(t) = \frac{-c[-ca(t) - cb(t)f(t)^{n-2}]}{(1 - cA(t) - c\int_{0}^{t}b(\tau)f(\tau)^{n-2}d\tau)^{2}}$$

$$= \frac{c^{2}a(t)}{(1 - cA(t) - c\int_{0}^{t}b(\tau)f(\tau)^{n-2}d\tau)^{2}} + \frac{c^{2}b(t)f(t)^{n-2}}{(1 - cA(t) - c\int_{0}^{t}b(\tau)f(\tau)^{n-2}d\tau)^{2}}$$

$$T_{c}(f)(t) - T_{c}(g)(t) = \frac{c}{H(f)} \cdot \frac{c}{H(g)} \cdot \int_{0}^{t}b(\tau)(f(\tau)^{n-2} - g(\tau)^{n-2})d\tau$$

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Lemma 2				

Lemma (2)

Let $c_1 = (||a|| + ||b|| + 1)^{-1}$. Then we have

 $||T_c f|| \le 1$ if $||f|| \le 1$ and $|c| \le c_1$.

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Lemma 2				

Lemma (2)

Let $c_1 = (||a|| + ||b|| + 1)^{-1}$. Then we have

$$||T_c f|| \le 1$$
 if $||f|| \le 1$ and $|c| \le c_1$.

Outline of the Proof.

If $||f|| \leq 1$ and $|c| \leq c_1$, then we have

$$\begin{aligned} \|T_{c}f\| &\leq \frac{|c|}{1-|c|\left(\|a\|+\|b\|\|f\|^{n-2}\right)} \\ &\leq \frac{|c|}{1-|c|\left(\|a\|+\|b\|\right)} \\ &< 1. \end{aligned}$$

The conclusion follows.

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Lemma 3				

Lemma (3)

Let $c_2 = (\sqrt{(n-2)||b||} + ||a|| + ||b|| + 1)^{-1}$. If $|c| \le c_2$, then T_c is a contraction mapping on the close unit ball $\mathcal{B}_1 = \{f \in \mathcal{C}[0,1] | ||f|| \le 1\}$ of $\mathcal{C}[0,1]$.

Outline of the Proof.

It follows from Lemmas 1 and 2 that

$$\begin{aligned} \|T_c(f)(t) - T_c(g)(t)\| &\leq \|T_c(f)\| \|T_c(g)\| \|b\| \|f^{n-2} - g^{n-2}\| \\ &= C\| (f-g)(f^{n-3} + f^{n-4}g + \dots + fg^{n-4} + g^{n-3})\| \\ &\leq (n-2)c\|f-g\|, \end{aligned}$$

where

$$c \stackrel{def}{=} \left(\frac{|c|}{1 - |c| \left(\|a\| + \|b\| \right)} \right)^2 \|b\|.$$

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Theorem (1)

For given $a, b \in C[0, 1]$ and $c \in \mathbb{R}$ with $|c| \leq (\sqrt{(n-2)}||b|| + ||a|| + ||b|| + 1)^{-1}$, the solution r(t, c) of equation (1) with r(0, c) = c can be uniformly approximated by an iterated sequence $\{T_c^n(f)(t)\}$:

$$r(t, c) = \lim_{n \to \infty} T_c^n(f)(t), \quad 0 \le t \le 1,$$
 (17)

that is,

$$r(t, c) = \frac{c}{1 - cA(t) - c^{n-1} \int_0^t \frac{b(t_1)dt_1}{1 - cA(t_1) - c^{n-1} \int_0^{t_1} \frac{b(t_2)dt_2}{1 - cA(t_2) - c^{n-1} \int_0^{t_2} \cdots}}$$
(18)

for arbitrary $f \in C[0, 1]$ with $||f|| \le 1$. Furthermore, the following error estimate holds

$$r(t, c) - T_c^n(f)(t) = \mathcal{O}(c^{2n}).$$

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Theorem 2	: Case of $n = 3$			

• Denote

$$M = \max_{t \in [0,1]} |a(t) \pm b(t)|.$$

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Theorem 2: (Case of $n = 3$			

• Denote

$$M = \max_{t \in [0,1]} |a(t) \pm b(t)|.$$

Theorem (2)

Suppose $a, b \in C[0, 1]$ and $c \in \mathbb{R}$ with

$$|c| \le \max\{(\sqrt{\|b\|} + \|a\| + \|b\| + 1)^{-1}, (2M)^{-1}\}.$$

Then, in formula (18), the following part is bounded

$$\begin{array}{rcl} & b(t_1) \\ \hline & \hline 1 - cA(t_1) - c^2 \int_0^{t_1} \frac{b(t_2)dt_2}{1 - cA(t_2) - c^2 \int_0^{t_2} \cdots} \\ & = & \frac{1}{c} \cdot b(t_1) \cdot \frac{c}{1 - cA(t_1) - c^2 \int_0^{t_1} b(t_2) \cdot \frac{c}{1 - cA(t_2) - c^2 \int_0^{t_2} \cdots} dt_2}. \end{array} \end{array}$$

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2D Korteweg-	de Vries-Burgers E	quation		

$$(U_t + \alpha UU_x + \beta U_{xx} + sU_{xxx})_x + \gamma U_{yy} = 0,$$
(19)

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2D Korteweg	-de Vries-Burgers H	Equation		

$$(U_t + \alpha UU_x + \beta U_{xx} + sU_{xxx})_x + \gamma U_{yy} = 0, \qquad (19)$$

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where α , β , *s*, and γ are constants and $\alpha\beta s\gamma \neq 0$.



$$(U_t + \alpha UU_x + \beta U_{xx} + sU_{xxx})_x + \gamma U_{yy} = 0, \qquad (19)$$

where α , β , s, and γ are constants and $\alpha\beta s\gamma \neq 0$.

• Assume that equation (19) has an exact solution in the form

$$U(x, y, t) = U(\xi), \quad \xi = hx + ly - wt.$$
 (20)

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$$(U_t + \alpha UU_x + \beta U_{xx} + sU_{xxx})_x + \gamma U_{yy} = 0, \qquad (19)$$

where α , β , s, and γ are constants and $\alpha\beta s\gamma \neq 0$.

• Assume that equation (19) has an exact solution in the form

$$U(x, y, t) = U(\xi), \quad \xi = hx + ly - wt.$$
 (20)

• Substitution of (20) into equation (19) and performing integration twice yields

$$U''(\xi) + \lambda U'(\xi) + aU^2(\xi) + bU(\xi) + d = 0,$$
(21)

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where $v = U(\xi) \in [v_0, v_1], \lambda = \frac{\beta}{sh}, a = \frac{\alpha}{2sh^2}, b = \frac{\gamma l^2 - wh}{sh^4}$ and $d = -\frac{C}{sh^4}$.

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Resultant /	Abel Equation			

• Let $v = U(\xi)$ and $y = U'(\xi)$. Equation (21) becomes

$$\frac{dy}{dv}y + \lambda y + av^2 + bv + d = 0.$$
(22)

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Resultant Abel Equation

• Let $v = U(\xi)$ and $y = U'(\xi)$. Equation (21) becomes

$$\frac{dy}{dv}y + \lambda y + av^2 + bv + d = 0.$$
(22)

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Using $z = \frac{1}{y}$ yields $\frac{dz}{dv} = \lambda z^2 + (av^2 + bv + d)z^3, \quad z(v_0) = \frac{1}{U'(\xi_0)} = c.$ (23)

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Resultant Abe	el Equation			

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$$\frac{dy}{dv}y + \lambda y + av^2 + bv + d = 0.$$
(22)

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Let $\eta = \frac{v - v_0}{v_1 - v_0}$, then $\eta \in [0, 1]$ and $v = v_0 + (v_1 - v_0)\eta$.

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Resultant Aber Equation

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• Let $\eta = \frac{v-v_0}{v_1-v_0}$, then $\eta \in [0, 1]$ and $v = v_0 + (v_1 - v_0)\eta$. So equation (23) reduces to

$$r' = h(\eta)r^2 + k(\eta)r^3, \quad r(0) = c,$$
 (24)

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Resultant Abel Equation

• Let $v = U(\xi)$ and $y = U'(\xi)$. Equation (21) becomes

$$\frac{dy}{dv}y + \lambda y + av^2 + bv + d = 0.$$
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• Let $\eta = \frac{v-v_0}{v_1-v_0}$, then $\eta \in [0, 1]$ and $v = v_0 + (v_1 - v_0)\eta$. So equation (23) reduces to

$$r' = h(\eta)r^2 + k(\eta)r^3, \quad r(0) = c,$$
 (24)

where $h(\eta), k(\eta) \in \mathcal{C}[0, 1]$, and

$$\begin{split} h(\eta) &= (v_1 - v_0)\lambda, \\ k(\eta) &= (v_1 - v_0)(av^2 + bv + d). \\ &\stackrel{\scriptstyle <}{\underset{\scriptstyle < \Box \, > \, < \, \Box \, > \, = \, \ldots \,$$

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Solution to	Equation (24)			
Solution to	Lqualion (2+)			

• By virtue of Theorem 1, if $|c| \le (\sqrt{\|k\|} + \|h\| + \|k\| + 1)^{-1}$, the solution to equation (24) is

$$r(\eta) = \lim_{n \to +\infty} T_c^n(w)(\eta), \tag{25}$$

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Solution to	Equation (24)			

• By virtue of Theorem 1, if $|c| \le (\sqrt{\|k\|} + \|h\| + \|k\| + 1)^{-1}$, the solution to equation (24) is

$$r(\eta) = \lim_{n \to +\infty} T_c^n(w)(\eta), \tag{25}$$

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where $0 \le \eta \le 1$ for any $w \in C[0, 1]$ with $||w|| \le 1$, and

$$T_{c}(w) = \frac{c}{1 - cH(\eta) - c\int_{0}^{\eta} k(x)w(x)^{n-2}dx}$$

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Solution to	Equation (24)			

• By virtue of Theorem 1, if $|c| \le (\sqrt{\|k\|} + \|h\| + \|k\| + 1)^{-1}$, the solution to equation (24) is

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where $0 \le \eta \le 1$ for any $w \in C[0, 1]$ with $||w|| \le 1$, and

$$T_{c}(w) = \frac{c}{1 - cH(\eta) - c \int_{0}^{\eta} k(x)w(x)^{n-2}dx}$$

where

$$H(\eta) = \int_0^{\eta} h(x) dx = \int_0^{\eta} (v_1 - v_0) \lambda dx = (v_1 - v_0) \lambda \eta,$$

$$k(x) = (v_1 - v_0) \left(a(v_0 + (v_1 - v_0)x)^2 + b(v_0 + (v_1 - v_0)x) + d \right).$$

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• Recall that $r = \frac{1}{y}$, $y = U'(\xi)$, $\eta = \frac{v - v_0}{v_1 - v_0}$ and $v = U(\xi)$. When conditions of Theorem 1 are fulfilled, we have

$$\frac{1}{U'(\xi)} = \frac{c}{1 - cA(\xi) - c^2 \int_0^{\xi} \frac{b(t_1)dt_1}{1 - cA(t_1) - c^2 \int_0^{t_1} \frac{b(t_2)dt_2}{1 - cA(t_2) - c^2 \int_0^{t_2} \dots}}.$$
 (26)

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• Recall that $r = \frac{1}{y}$, $y = U'(\xi)$, $\eta = \frac{v - v_0}{v_1 - v_0}$ and $v = U(\xi)$. When conditions of Theorem 1 are fulfilled, we have

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 (26)

• When c is small, according to Theorem 2, the coefficient of c^2 is bounded. So we can drop the term containing c^2 and get

$$U'(\xi) \approx \frac{1 - c(v_1 - v_0)\lambda\eta}{c}$$
$$= \frac{1 - c\lambda(U(\xi) - v_0)}{c}$$

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• Recall that $r = \frac{1}{y}$, $y = U'(\xi)$, $\eta = \frac{v - v_0}{v_1 - v_0}$ and $v = U(\xi)$. When conditions of Theorem 1 are fulfilled, we have

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 (26)

• When *c* is small, according to Theorem 2, the coefficient of *c*² is bounded. So we can drop the term containing *c*² and get

$$U^{'}(\xi) \approx \frac{1 - c(v_1 - v_0)\lambda\eta}{c}$$
$$= \frac{1 - c\lambda(U(\xi) - v_0)}{c}$$

That is,

$$U'(\xi) + \lambda U(\xi) = \frac{1}{c} + \lambda v_0. \tag{27}$$

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• Solving equation (27) gives

$$U(x, y, t) = \frac{\frac{1}{c} + \lambda v_0}{\lambda} + ce^{-\lambda\xi}, \quad \xi = hx + ly - wt$$

where $\lambda = \frac{\beta}{sh}$.

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• Solving equation (27) gives

$$U(x, y, t) = \frac{\frac{1}{c} + \lambda v_0}{\lambda} + ce^{-\lambda\xi}, \quad \xi = hx + ly - wt$$

where $\lambda = \frac{\beta}{sh}$.

• If we take $v_0 = \frac{b}{2a}$ and choose $c = \frac{-2a}{\lambda\sqrt{b^2-4ad}}$ sufficiently small, when $\lambda\xi \to +\infty$, we obtain

$$U(x, y, t) \sim \frac{b^2 - 4ad}{-2a} + \frac{b}{2a}.$$
 (28)

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• Solving equation (27) gives

$$U(x, y, t) = \frac{\frac{1}{c} + \lambda v_0}{\lambda} + ce^{-\lambda\xi}, \quad \xi = hx + ly - wt$$

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 (28)

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• It is remarkable that the approximate solution (28) is in agreement with main results described in [7, 8] by the Hardy's theory and the theory of Lie symmetry.

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• Solving equation (27) gives

$$U(x, y, t) = \frac{\frac{1}{c} + \lambda v_0}{\lambda} + ce^{-\lambda\xi}, \quad \xi = hx + ly - wt$$

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• It is remarkable that the approximate solution (28) is in agreement with main results described in [7, 8] by the Hardy's theory and the theory of Lie symmetry.

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• Solving equation (27) gives

$$U(x, y, t) = \frac{\frac{1}{c} + \lambda v_0}{\lambda} + ce^{-\lambda\xi}, \quad \xi = hx + ly - wt$$

where $\lambda = \frac{\beta}{sh}$.

• If we take $v_0 = \frac{b}{2a}$ and choose $c = \frac{-2a}{\lambda\sqrt{b^2-4ad}}$ sufficiently small, when $\lambda\xi \to +\infty$, we obtain

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[8] Z. Feng, Nonlinearity, 20 (2007), 343–356.

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Boundedness	of Solutions			

• Note that equation (26) can be rewritten as

$$\frac{1}{U'(\xi)} = \frac{c}{1 - cA(\xi) - c^2 \Phi(\xi)},$$
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where $L \le \Phi(\xi) \le R$.

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where $L \leq \Phi(\xi) \leq R$.

• When Φ is a quadratic or cubic function or special function of $U(\xi)$, we can analyze equation (29) qualitatively and numerically with classifications. For instance, if Φ is quadratic, we take $v_0 = \frac{b}{2a}$ and choose $c = \frac{-2a}{\lambda\sqrt{b^2-4ad}}$ sufficiently small, we can obtain the solution of the type

$$u(x, y, t) = \frac{3\beta^2 + \gamma + c}{25\alpha s} \operatorname{sech}^2 \xi - \frac{6\beta^2 + \gamma + c}{25\alpha s} \tanh \xi \pm \frac{6\beta^2}{25\alpha s} + C_0.$$

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 When Φ is a function with the lower and upper bounds, we can also find bounds of solutions of equation (29) by the comparison principle, which match well with the phase analysis described in [7].

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Summary				

• In this talk, we provided a connection between the Abel equation of the first kind, an ordinary differential equation that is cubic in the unknown function, and the Korteweg-de Vries-Burgers equation, a partial differential equation that describes the propagation of waves on liquid-filled elastic tubes. We presented an integral form of the Abel equation with the initial condition.

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- By virtue of the integral form and the Banach Contraction Mapping Principle we derived the asymptotic expansion of bounded solutions in the Banach space, and used the asymptotic formula to construct approximate solutions to the Korteweg-de Vries-Burgers equation.

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- As an example, we presented the asymptotic behavior of traveling wave solution for a 2D KdV-Burgers equation which agrees well with existing results in the literature.
- Under certain conditions, we can also study bounds of traveling wave solutions of KdV-Burgers type equations by the comparison principle.

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• I would like to thank

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• I would like to thank Xiaoqian Gong for discussions and help on computations.

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- I would like to thank Xiaoqian Gong for discussions and help on computations.
- Thank you.

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