# Approximate Analysis to the KdV-Burgers Equation 

Zhaosheng Feng

Department of Mathematics
University of Texas-Pan American
1201 W. University Dr.
Edinburg, Texas 78539, USA
E-mail: zsfeng@utpa.edu

October 26, 2013

Texas Analysis and Mathematical Physics Symposium—Rice University

## Outline

## (1) Introduction

- Generalized KdV-Burgers Equation
- KdV-Burgers Equation
- Planar Polynomial Systems and Abel Equation


## Outline

(1) Introduction

- Generalized KdV-Burgers Equation
- KdV-Burgers Equation
- Planar Polynomial Systems and Abel Equation
(2) Qualitative Analysis
- Generalized Abel Equation
- Property of Our Operator
- Two Theorems


## Outline

(1) Introduction

- Generalized KdV-Burgers Equation
- KdV-Burgers Equation
- Planar Polynomial Systems and Abel Equation
(2) Qualitative Analysis
- Generalized Abel Equation
- Property of Our Operator
- Two Theorems
(3) Approximate Solution
- 2D KdV-Burgers Equation
- Resultant Abel Equation
- Approximate Solution to 2D KdV-Burgers Equation


## Outline

(1) Introduction

- Generalized KdV-Burgers Equation
- KdV-Burgers Equation
- Planar Polynomial Systems and Abel Equation
(2) Qualitative Analysis
- Generalized Abel Equation
- Property of Our Operator
- Two Theorems
(3) Approximate Solution
- 2D KdV-Burgers Equation
- Resultant Abel Equation
- Approximate Solution to 2D KdV-Burgers Equation
(4) Conclusion


## Outline

(1) Introduction

- Generalized KdV-Burgers Equation
- KdV-Burgers Equation
- Planar Polynomial Systems and Abel Equation
(2) Qualitative Analysis
- Generalized Abel Equation
- Property of Our Operator
- Two Theorems
(3) Approximate Solution
- 2D KdV-Burgers Equation
- Resultant Abel Equation
- Approximate Solution to 2D KdV-Burgers Equation
(4) Conclusion
(5) Acknowledgement


## Generalized KdV-Burgers Equation

- Generalized Korteweg-de Vries-Burgers equation [1, 2]

$$
\begin{equation*}
u_{t}+\left(\delta u_{x x}+\frac{\beta}{p} u^{p}\right)_{x}+\alpha u_{x}-\mu u_{x x}=0 \tag{1}
\end{equation*}
$$

where $u$ is a function of $x$ and $t, \alpha, \beta$ and $p>0$ are real constants, $\mu$ and $\delta$ are coefficients of dissipation and dispersion, respectively.

## Generalized KdV-Burgers Equation

- Generalized Korteweg-de Vries-Burgers equation [1, 2]

$$
\begin{equation*}
u_{t}+\left(\delta u_{x x}+\frac{\beta}{p} u^{p}\right)_{x}+\alpha u_{x}-\mu u_{x x}=0 \tag{1}
\end{equation*}
$$

where $u$ is a function of $x$ and $t, \alpha, \beta$ and $p>0$ are real constants, $\mu$ and $\delta$ are coefficients of dissipation and dispersion, respectively.

- The type of such problems arises in modeling waves generated by a wavemaker in a channel and waves incoming from deep water into nearshore zones.


## Generalized KdV-Burgers Equation

- Generalized Korteweg-de Vries-Burgers equation [1, 2]

$$
\begin{equation*}
u_{t}+\left(\delta u_{x x}+\frac{\beta}{p} u^{p}\right)_{x}+\alpha u_{x}-\mu u_{x x}=0 \tag{1}
\end{equation*}
$$

where $u$ is a function of $x$ and $t, \alpha, \beta$ and $p>0$ are real constants, $\mu$ and $\delta$ are coefficients of dissipation and dispersion, respectively.

- The type of such problems arises in modeling waves generated by a wavemaker in a channel and waves incoming from deep water into nearshore zones.
- The type of such models has the simplest form of wave equation in which nonlinearity $\left(u^{p}\right)_{x}$, dispersion $u_{x x x}$ and dissipation $u_{x x}$ all occur.


## Generalized KdV-Burgers Equation

- Generalized Korteweg-de Vries-Burgers equation [1, 2]

$$
\begin{equation*}
u_{t}+\left(\delta u_{x x}+\frac{\beta}{p} u^{p}\right)_{x}+\alpha u_{x}-\mu u_{x x}=0 \tag{1}
\end{equation*}
$$

where $u$ is a function of $x$ and $t, \alpha, \beta$ and $p>0$ are real constants, $\mu$ and $\delta$ are coefficients of dissipation and dispersion, respectively.

- The type of such problems arises in modeling waves generated by a wavemaker in a channel and waves incoming from deep water into nearshore zones.
- The type of such models has the simplest form of wave equation in which nonlinearity $\left(u^{p}\right)_{x}$, dispersion $u_{x x x}$ and dissipation $u_{x x}$ all occur.


## Generalized KdV-Burgers Equation

- Generalized Korteweg-de Vries-Burgers equation [1, 2]

$$
\begin{equation*}
u_{t}+\left(\delta u_{x x}+\frac{\beta}{p} u^{p}\right)_{x}+\alpha u_{x}-\mu u_{x x}=0 \tag{1}
\end{equation*}
$$

where $u$ is a function of $x$ and $t, \alpha, \beta$ and $p>0$ are real constants, $\mu$ and $\delta$ are coefficients of dissipation and dispersion, respectively.

- The type of such problems arises in modeling waves generated by a wavemaker in a channel and waves incoming from deep water into nearshore zones.
- The type of such models has the simplest form of wave equation in which nonlinearity $\left(u^{p}\right)_{x}$, dispersion $u_{x x x}$ and dissipation $u_{x x}$ all occur.
[1] J.L. Bona, W.G. Pritchard and L.R. Scott, Philos. Trans. Roy. Soc. London Ser. A, 302 (1981), 457-510.


## Generalized KdV-Burgers Equation

- Generalized Korteweg-de Vries-Burgers equation [1, 2]

$$
\begin{equation*}
u_{t}+\left(\delta u_{x x}+\frac{\beta}{p} u^{p}\right)_{x}+\alpha u_{x}-\mu u_{x x}=0 \tag{1}
\end{equation*}
$$

where $u$ is a function of $x$ and $t, \alpha, \beta$ and $p>0$ are real constants, $\mu$ and $\delta$ are coefficients of dissipation and dispersion, respectively.

- The type of such problems arises in modeling waves generated by a wavemaker in a channel and waves incoming from deep water into nearshore zones.
- The type of such models has the simplest form of wave equation in which nonlinearity $\left(u^{p}\right)_{x}$, dispersion $u_{x x x}$ and dissipation $u_{x x}$ all occur.
[1] J.L. Bona, W.G. Pritchard and L.R. Scott, Philos. Trans. Roy. Soc. London Ser. A, 302 (1981), 457-510.
[2] J.L.Bona, S.M. Sun and B.Y. Zhang, Dyn. Partial Differ. Equs. 3 (2006), 1-69.


## Burgers Equation and KdV Equation

- Choices of $\delta=\alpha=0$ and $p=2$ lead (1) to the Burgers equation [3]:

$$
\begin{equation*}
u_{t}+\alpha u u_{x}+\beta u_{x x}=0, \tag{2}
\end{equation*}
$$

## Burgers Equation and KdV Equation

- Choices of $\delta=\alpha=0$ and $p=2$ lead (1) to the Burgers equation [3]:

$$
\begin{equation*}
u_{t}+\alpha u u_{x}+\beta u_{x x}=0, \tag{2}
\end{equation*}
$$

with the wave solution

$$
u(x, t)=\frac{2 k}{\alpha}+\frac{2 \beta k}{\alpha} \tanh k(x-2 k t) .
$$

## Burgers Equation and KdV Equation

- Choices of $\delta=\alpha=0$ and $p=2$ lead (1) to the Burgers equation [3]:

$$
\begin{equation*}
u_{t}+\alpha u u_{x}+\beta u_{x x}=0, \tag{2}
\end{equation*}
$$

with the wave solution

$$
u(x, t)=\frac{2 k}{\alpha}+\frac{2 \beta k}{\alpha} \tanh k(x-2 k t) .
$$

- Choices of $\alpha=\mu=0$ and $p=2$ lead (1) to the KdV equation [4]:

$$
\begin{equation*}
u_{t}+\alpha u u_{x}+s u_{x x x}=0, \tag{3}
\end{equation*}
$$

## Burgers Equation and KdV Equation

- Choices of $\delta=\alpha=0$ and $p=2$ lead (1) to the Burgers equation [3]:

$$
\begin{equation*}
u_{t}+\alpha u u_{x}+\beta u_{x x}=0, \tag{2}
\end{equation*}
$$

with the wave solution

$$
u(x, t)=\frac{2 k}{\alpha}+\frac{2 \beta k}{\alpha} \tanh k(x-2 k t) .
$$

- Choices of $\alpha=\mu=0$ and $p=2$ lead (1) to the KdV equation [4]:

$$
\begin{equation*}
u_{t}+\alpha u u_{x}+s u_{x x x}=0, \tag{3}
\end{equation*}
$$

with the soliton solution [5]

$$
u(x, t)=\frac{12 s k^{2}}{\alpha} \operatorname{sech}^{2} k\left(x-4 s k^{2} t\right)
$$

## Burgers Equation and KdV Equation

- Choices of $\delta=\alpha=0$ and $p=2$ lead (1) to the Burgers equation [3]:

$$
\begin{equation*}
u_{t}+\alpha u u_{x}+\beta u_{x x}=0, \tag{2}
\end{equation*}
$$

with the wave solution

$$
u(x, t)=\frac{2 k}{\alpha}+\frac{2 \beta k}{\alpha} \tanh k(x-2 k t) .
$$

- Choices of $\alpha=\mu=0$ and $p=2$ lead (1) to the KdV equation [4]:

$$
\begin{equation*}
u_{t}+\alpha u u_{x}+s u_{x x x}=0, \tag{3}
\end{equation*}
$$

with the soliton solution [5]

$$
u(x, t)=\frac{12 s k^{2}}{\alpha} \operatorname{sech}^{2} k\left(x-4 s k^{2} t\right)
$$

[3] J.M. Burgers, Trans. Roy. Neth. Acad. Sci. 17 (1939), 1-53

## Burgers Equation and KdV Equation

- Choices of $\delta=\alpha=0$ and $p=2$ lead (1) to the Burgers equation [3]:

$$
\begin{equation*}
u_{t}+\alpha u u_{x}+\beta u_{x x}=0, \tag{2}
\end{equation*}
$$

with the wave solution

$$
u(x, t)=\frac{2 k}{\alpha}+\frac{2 \beta k}{\alpha} \tanh k(x-2 k t) .
$$

- Choices of $\alpha=\mu=0$ and $p=2$ lead (1) to the KdV equation [4]:

$$
\begin{equation*}
u_{t}+\alpha u u_{x}+s u_{x x x}=0, \tag{3}
\end{equation*}
$$

with the soliton solution [5]

$$
u(x, t)=\frac{12 s k^{2}}{\alpha} \operatorname{sech}^{2} k\left(x-4 s k^{2} t\right)
$$

[3] J.M. Burgers, Trans. Roy. Neth. Acad. Sci. 17 (1939), 1-53
[4] D.J. Korteweg and G. de Vries, Phil. Mag. 39 (1895), 422-443.

## Burgers Equation and KdV Equation

- Choices of $\delta=\alpha=0$ and $p=2$ lead (1) to the Burgers equation [3]:

$$
\begin{equation*}
u_{t}+\alpha u u_{x}+\beta u_{x x}=0, \tag{2}
\end{equation*}
$$

with the wave solution

$$
u(x, t)=\frac{2 k}{\alpha}+\frac{2 \beta k}{\alpha} \tanh k(x-2 k t) .
$$

- Choices of $\alpha=\mu=0$ and $p=2$ lead (1) to the KdV equation [4]:

$$
\begin{equation*}
u_{t}+\alpha u u_{x}+s u_{x x x}=0, \tag{3}
\end{equation*}
$$

with the soliton solution [5]

$$
u(x, t)=\frac{12 s k^{2}}{\alpha} \operatorname{sech}^{2} k\left(x-4 s k^{2} t\right)
$$

[3] J.M. Burgers, Trans. Roy. Neth. Acad. Sci. 17 (1939), 1-53
[4] D.J. Korteweg and G. de Vries, Phil. Mag. 39 (1895), 422-443.
[5] N.J. Zabusky and M.D. Kruskal, Phys. Rev. Lett. 15 (1965), 240-243.

## Korteweg-de Vries-Burgers Equation

- Choices of $\alpha=0$ and $p=2$ lead equation (1) to the standard form of the Korteweg-de Vries-Burgers equation [6]:

$$
\begin{equation*}
u_{t}+\alpha u u_{x}+\beta u_{x x}+s u_{x x x}=0 \tag{4}
\end{equation*}
$$

## Korteweg-de Vries-Burgers Equation

- Choices of $\alpha=0$ and $p=2$ lead equation (1) to the standard form of the Korteweg-de Vries-Burgers equation [6]:

$$
\begin{equation*}
u_{t}+\alpha u u_{x}+\beta u_{x x}+s u_{x x x}=0 . \tag{4}
\end{equation*}
$$

- Solitary wave solutions of equation (4) are as follows [7, 8, 9]:

$$
\begin{equation*}
u(x, t)=\frac{3 \beta^{2}}{25 \alpha s} \operatorname{sech}^{2} \Psi-\frac{6 \beta^{2}}{25 \alpha s} \tanh \Psi \pm \frac{6 \beta^{2}}{25 \alpha s} \tag{5}
\end{equation*}
$$

## Korteweg-de Vries-Burgers Equation

- Choices of $\alpha=0$ and $p=2$ lead equation (1) to the standard form of the Korteweg-de Vries-Burgers equation [6]:

$$
\begin{equation*}
u_{t}+\alpha u u_{x}+\beta u_{x x}+s u_{x x x}=0 . \tag{4}
\end{equation*}
$$

- Solitary wave solutions of equation (4) are as follows [7, 8, 9]:

$$
\begin{equation*}
u(x, t)=\frac{3 \beta^{2}}{25 \alpha s} \operatorname{sech}^{2} \Psi-\frac{6 \beta^{2}}{25 \alpha s} \tanh \Psi \pm \frac{6 \beta^{2}}{25 \alpha s} \tag{5}
\end{equation*}
$$

where

$$
\Psi=\left[\frac{1}{2}\left(-\frac{\beta}{5 s} x \pm \frac{6 \beta^{3}}{125 s^{2}} t\right)\right]
$$

## Korteweg-de Vries-Burgers Equation

- Choices of $\alpha=0$ and $p=2$ lead equation (1) to the standard form of the Korteweg-de Vries-Burgers equation [6]:

$$
\begin{equation*}
u_{t}+\alpha u u_{x}+\beta u_{x x}+s u_{x x x}=0 . \tag{4}
\end{equation*}
$$

- Solitary wave solutions of equation (4) are as follows [7, 8, 9]:

$$
\begin{equation*}
u(x, t)=\frac{3 \beta^{2}}{25 \alpha s} \operatorname{sech}^{2} \Psi-\frac{6 \beta^{2}}{25 \alpha s} \tanh \Psi \pm \frac{6 \beta^{2}}{25 \alpha s} \tag{5}
\end{equation*}
$$

where

$$
\Psi=\left[\frac{1}{2}\left(-\frac{\beta}{5 s} x \pm \frac{6 \beta^{3}}{125 s^{2}} t\right)\right]
$$

## Korteweg-de Vries-Burgers Equation

- Choices of $\alpha=0$ and $p=2$ lead equation (1) to the standard form of the Korteweg-de Vries-Burgers equation [6]:

$$
\begin{equation*}
u_{t}+\alpha u u_{x}+\beta u_{x x}+s u_{x x x}=0 . \tag{4}
\end{equation*}
$$

- Solitary wave solutions of equation (4) are as follows [7, 8, 9]:

$$
\begin{equation*}
u(x, t)=\frac{3 \beta^{2}}{25 \alpha s} \operatorname{sech}^{2} \Psi-\frac{6 \beta^{2}}{25 \alpha s} \tanh \Psi \pm \frac{6 \beta^{2}}{25 \alpha s} \tag{5}
\end{equation*}
$$

where

$$
\Psi=\left[\frac{1}{2}\left(-\frac{\beta}{5 s} x \pm \frac{6 \beta^{3}}{125 s^{2}} t\right)\right] .
$$

[6] R.S. Johnson, J. Fluid Mech. 42 (1970), 49-60.

## Korteweg-de Vries-Burgers Equation

- Choices of $\alpha=0$ and $p=2$ lead equation (1) to the standard form of the Korteweg-de Vries-Burgers equation [6]:

$$
\begin{equation*}
u_{t}+\alpha u u_{x}+\beta u_{x x}+s u_{x x x}=0 . \tag{4}
\end{equation*}
$$

- Solitary wave solutions of equation (4) are as follows [7, 8, 9]:

$$
\begin{equation*}
u(x, t)=\frac{3 \beta^{2}}{25 \alpha s} \operatorname{sech}^{2} \Psi-\frac{6 \beta^{2}}{25 \alpha s} \tanh \Psi \pm \frac{6 \beta^{2}}{25 \alpha s} \tag{5}
\end{equation*}
$$

where

$$
\Psi=\left[\frac{1}{2}\left(-\frac{\beta}{5 s} x \pm \frac{6 \beta^{3}}{125 s^{2}} t\right)\right] .
$$

[6] R.S. Johnson, J. Fluid Mech. 42 (1970), 49-60.
[7] Z. Feng, J. Phys. A (Math. Gen.) 36 (2003), 8817-8827.

## Korteweg-de Vries-Burgers Equation

- Choices of $\alpha=0$ and $p=2$ lead equation (1) to the standard form of the Korteweg-de Vries-Burgers equation [6]:

$$
\begin{equation*}
u_{t}+\alpha u u_{x}+\beta u_{x x}+s u_{x x x}=0 . \tag{4}
\end{equation*}
$$

- Solitary wave solutions of equation (4) are as follows [7, 8, 9]:

$$
\begin{equation*}
u(x, t)=\frac{3 \beta^{2}}{25 \alpha s} \operatorname{sech}^{2} \Psi-\frac{6 \beta^{2}}{25 \alpha s} \tanh \Psi \pm \frac{6 \beta^{2}}{25 \alpha s} \tag{5}
\end{equation*}
$$

where

$$
\Psi=\left[\frac{1}{2}\left(-\frac{\beta}{5 s} x \pm \frac{6 \beta^{3}}{125 s^{2}} t\right)\right] .
$$

[6] R.S. Johnson, J. Fluid Mech. 42 (1970), 49-60.
[7] Z. Feng, J. Phys. A (Math. Gen.) 36 (2003), 8817-8827.
[8] Z. Feng, Nonlinearity, 20 (2007), 343-356.

## Korteweg-de Vries-Burgers Equation

- Choices of $\alpha=0$ and $p=2$ lead equation (1) to the standard form of the Korteweg-de Vries-Burgers equation [6]:

$$
\begin{equation*}
u_{t}+\alpha u u_{x}+\beta u_{x x}+s u_{x x x}=0 . \tag{4}
\end{equation*}
$$

- Solitary wave solutions of equation (4) are as follows [7, 8, 9]:

$$
\begin{equation*}
u(x, t)=\frac{3 \beta^{2}}{25 \alpha s} \operatorname{sech}^{2} \Psi-\frac{6 \beta^{2}}{25 \alpha s} \tanh \Psi \pm \frac{6 \beta^{2}}{25 \alpha s}, \tag{5}
\end{equation*}
$$

where

$$
\Psi=\left[\frac{1}{2}\left(-\frac{\beta}{5 s} x \pm \frac{6 \beta^{3}}{125 s^{2}} t\right)\right] .
$$

[6] R.S. Johnson, J. Fluid Mech. 42 (1970), 49-60.
[7] Z. Feng, J. Phys. A (Math. Gen.) 36 (2003), 8817-8827.
[8] Z. Feng, Nonlinearity, 20 (2007), 343-356.
[9] Z. Feng and S. Zheng, Z. angew. Math. Phys. 60 (2009), 756-773.

## Figures of Wave Solutions



Legend

- u1-Burgers-KdV u2-Burgers-KdV u3-Burgers u4-KdV


## Planar Polynomial Systems and Abel Equation

- Consider planar polynomial systems of the form

$$
\begin{equation*}
\dot{x}=-y+p(x, y), \quad \dot{y}=x+q(x, y) \tag{6}
\end{equation*}
$$

## Planar Polynomial Systems and Abel Equation

- Consider planar polynomial systems of the form

$$
\begin{equation*}
\dot{x}=-y+p(x, y), \quad \dot{y}=x+q(x, y) \tag{6}
\end{equation*}
$$

with homogeneous polynomials $p(x, y)$ and $q(x, y)$ of degree $k$.

## Planar Polynomial Systems and Abel Equation

- Consider planar polynomial systems of the form

$$
\begin{equation*}
\dot{x}=-y+p(x, y), \quad \dot{y}=x+q(x, y) \tag{6}
\end{equation*}
$$

with homogeneous polynomials $p(x, y)$ and $q(x, y)$ of degree $k$.

- For the Poincaré center problem, setting $x=r \cos \theta, y=r \sin \theta$ gives

$$
\begin{equation*}
\frac{d r}{d \theta}=\frac{r^{k} \xi(\theta)}{1+r^{k-1} \eta(\theta)} \tag{7}
\end{equation*}
$$

where $\xi$ and $\eta$ are polynomials in $\cos \theta$ and $\sin \theta$ of degree $k+1$.

## Planar Polynomial Systems and Abel Equation

- Consider planar polynomial systems of the form

$$
\begin{equation*}
\dot{x}=-y+p(x, y), \quad \dot{y}=x+q(x, y) \tag{6}
\end{equation*}
$$

with homogeneous polynomials $p(x, y)$ and $q(x, y)$ of degree $k$.

- For the Poincaré center problem, setting $x=r \cos \theta, y=r \sin \theta$ gives

$$
\begin{equation*}
\frac{d r}{d \theta}=\frac{r^{k} \xi(\theta)}{1+r^{k-1} \eta(\theta)} \tag{7}
\end{equation*}
$$

where $\xi$ and $\eta$ are polynomials in $\cos \theta$ and $\sin \theta$ of degree $k+1$.

- Making the coordinate transformation

$$
\rho=\frac{r^{k-1}}{1+r^{k-1} \eta(\theta)},
$$

## Planar Polynomial Systems and Abel Equation

- Consider planar polynomial systems of the form

$$
\begin{equation*}
\dot{x}=-y+p(x, y), \quad \dot{y}=x+q(x, y) \tag{6}
\end{equation*}
$$

with homogeneous polynomials $p(x, y)$ and $q(x, y)$ of degree $k$.

- For the Poincaré center problem, setting $x=r \cos \theta, y=r \sin \theta$ gives

$$
\begin{equation*}
\frac{d r}{d \theta}=\frac{r^{k} \xi(\theta)}{1+r^{k-1} \eta(\theta)} \tag{7}
\end{equation*}
$$

where $\xi$ and $\eta$ are polynomials in $\cos \theta$ and $\sin \theta$ of degree $k+1$.

- Making the coordinate transformation

$$
\rho=\frac{r^{k-1}}{1+r^{k-1} \eta(\theta)},
$$

we get an Abel equation

$$
\frac{d \rho}{d \theta}=a(\theta) \rho^{2}+b(\theta) \rho^{3}
$$

where $a=(k-1) \xi+n^{\prime}$ and $b=(1-k) \xi \eta$.

## Traveling Wave Solution

- Assume that equation (1) has the traveling wave solution of the form

$$
u(x, t)=u(\xi), \quad \xi=x-c t,
$$

## Traveling Wave Solution

- Assume that equation (1) has the traveling wave solution of the form

$$
u(x, t)=u(\xi), \quad \xi=x-c t,
$$

where $c \neq 0$ is the wave velocity. Then equation (1) becomes

$$
\begin{equation*}
\delta u^{\prime \prime \prime}-\mu u^{\prime \prime}+(\alpha-c) u^{\prime}+\beta u^{p-1} u^{\prime}=0, \tag{8}
\end{equation*}
$$

## Traveling Wave Solution

- Assume that equation (1) has the traveling wave solution of the form

$$
u(x, t)=u(\xi), \quad \xi=x-c t,
$$

where $c \neq 0$ is the wave velocity. Then equation (1) becomes

$$
\begin{equation*}
\delta u^{\prime \prime \prime}-\mu u^{\prime \prime}+(\alpha-c) u^{\prime}+\beta u^{p-1} u^{\prime}=0, \tag{8}
\end{equation*}
$$

where $u^{\prime}=d u / d \xi$. Integrating equation (8) once gives

$$
\begin{equation*}
u^{\prime \prime}-g u^{\prime}-e u-f u^{p}-d=0, \tag{9}
\end{equation*}
$$

## Traveling Wave Solution

- Assume that equation (1) has the traveling wave solution of the form

$$
u(x, t)=u(\xi), \quad \xi=x-c t,
$$

where $c \neq 0$ is the wave velocity. Then equation (1) becomes

$$
\begin{equation*}
\delta u^{\prime \prime \prime}-\mu u^{\prime \prime}+(\alpha-c) u^{\prime}+\beta u^{p-1} u^{\prime}=0, \tag{8}
\end{equation*}
$$

where $u^{\prime}=d u / d \xi$. Integrating equation (8) once gives

$$
\begin{equation*}
u^{\prime \prime}-g u^{\prime}-e u-f u^{p}-d=0, \tag{9}
\end{equation*}
$$

where $e=\frac{c-\alpha}{\delta}, g=\frac{\mu}{\delta}, f=-\frac{\beta}{p \delta}$ and $d$ is an integration constant.

## Traveling Wave Solution

- Assume that equation (1) has the traveling wave solution of the form

$$
u(x, t)=u(\xi), \quad \xi=x-c t,
$$

where $c \neq 0$ is the wave velocity. Then equation (1) becomes

$$
\begin{equation*}
\delta u^{\prime \prime \prime}-\mu u^{\prime \prime}+(\alpha-c) u^{\prime}+\beta u^{p-1} u^{\prime}=0, \tag{8}
\end{equation*}
$$

where $u^{\prime}=d u / d \xi$. Integrating equation (8) once gives

$$
\begin{equation*}
u^{\prime \prime}-g u^{\prime}-e u-f u^{p}-d=0, \tag{9}
\end{equation*}
$$

where $e=\frac{c-\alpha}{\delta}, g=\frac{\mu}{\delta}, f=-\frac{\beta}{p \delta}$ and $d$ is an integration constant.

- Assume that $y=u$ and $u^{\prime}=z$, then equation (9) is equivalent to

$$
\left\{\begin{array}{l}
y^{\prime}=z  \tag{10}\\
z^{\prime}=e y+g z+f y^{p}+d .
\end{array}\right.
$$

## Global Structure of $p=2$



## Transformed to Abel Equation

- It follows from system (10) that

$$
\begin{equation*}
\frac{d z}{d y}=\frac{e y+g z+f y^{p}+d}{z} . \tag{11}
\end{equation*}
$$

## Transformed to Abel Equation

- It follows from system (10) that

$$
\begin{equation*}
\frac{d z}{d y}=\frac{e y+g z+f y^{p}+d}{z} \tag{11}
\end{equation*}
$$

- Let $z=r^{-1}$. Equation (11) reduces to

$$
\begin{equation*}
\frac{d r}{d y}=a(y) r^{2}+b(y) r^{3} \tag{12}
\end{equation*}
$$

## Transformed to Abel Equation

- It follows from system (10) that

$$
\begin{equation*}
\frac{d z}{d y}=\frac{e y+g z+f y^{p}+d}{z} \tag{11}
\end{equation*}
$$

- Let $z=r^{-1}$. Equation (11) reduces to

$$
\begin{equation*}
\frac{d r}{d y}=a(y) r^{2}+b(y) r^{3} \tag{12}
\end{equation*}
$$

where $a(y)=-g$ and $b(y)=-\left(e y+f y^{p}+d\right)$.

## Transformed to Abel Equation

- It follows from system (10) that

$$
\begin{equation*}
\frac{d z}{d y}=\frac{e y+g z+f y^{p}+d}{z} \tag{11}
\end{equation*}
$$

- Let $z=r^{-1}$. Equation (11) reduces to

$$
\begin{equation*}
\frac{d r}{d y}=a(y) r^{2}+b(y) r^{3} \tag{12}
\end{equation*}
$$

where $a(y)=-g$ and $b(y)=-\left(e y+f y^{p}+d\right)$.

- Question: Under what condition one can determine the number of closed solutions of the Abel equation (12).


## Transformed to Abel Equation

- It follows from system (10) that

$$
\begin{equation*}
\frac{d z}{d y}=\frac{e y+g z+f y^{p}+d}{z} \tag{11}
\end{equation*}
$$

- Let $z=r^{-1}$. Equation (11) reduces to

$$
\begin{equation*}
\frac{d r}{d y}=a(y) r^{2}+b(y) r^{3} \tag{12}
\end{equation*}
$$

where $a(y)=-g$ and $b(y)=-\left(e y+f y^{p}+d\right)$.

- Question: Under what condition one can determine the number of closed solutions of the Abel equation (12).
- Open Problem: There have been two longstanding problems, called the Poincaré center-focus problem and the local Hilbert 16th problem. Both are closely related to the Bautin quantities and the Bautin ideal of the Abel equation.


## Integral Form

- Consider the generalized Abel equation

$$
\begin{equation*}
r^{\prime}=a(t) r^{2}+b(t) r^{n}, \quad r\left(t_{0}\right)=c, \quad t \in\left[t_{0}, t_{1}\right], \quad n \geq 3 . \tag{13}
\end{equation*}
$$

## Integral Form

- Consider the generalized Abel equation

$$
\begin{equation*}
r^{\prime}=a(t) r^{2}+b(t) r^{n}, \quad r\left(t_{0}\right)=c, \quad t \in\left[t_{0}, t_{1}\right], \quad n \geq 3 . \tag{13}
\end{equation*}
$$

- Dividing both sides of equation (13) by $r^{2}$ gives

$$
\begin{equation*}
\frac{r^{\prime}}{r^{2}}=a(t)+b(t) r^{n-2} \tag{14}
\end{equation*}
$$

## Integral Form

- Consider the generalized Abel equation

$$
\begin{equation*}
r^{\prime}=a(t) r^{2}+b(t) r^{n}, \quad r\left(t_{0}\right)=c, \quad t \in\left[t_{0}, t_{1}\right], \quad n \geq 3 \tag{13}
\end{equation*}
$$

- Dividing both sides of equation (13) by $r^{2}$ gives

$$
\begin{equation*}
\frac{r^{\prime}}{r^{2}}=a(t)+b(t) r^{n-2} \tag{14}
\end{equation*}
$$

- Integrating equation (14) from $t_{0}$ to $t$ yields

$$
\begin{equation*}
r(t)=\frac{c}{1-c A(t)-c \int_{t_{0}}^{t} b(\tau) r^{n-2} d \tau} \tag{15}
\end{equation*}
$$

where $A(t)=\int_{t_{0}}^{t} a(\tau) d \tau$.

## Integral Form

- Consider the generalized Abel equation

$$
\begin{equation*}
r^{\prime}=a(t) r^{2}+b(t) r^{n}, \quad r\left(t_{0}\right)=c, \quad t \in\left[t_{0}, t_{1}\right], \quad n \geq 3 \tag{13}
\end{equation*}
$$

- Dividing both sides of equation (13) by $r^{2}$ gives

$$
\begin{equation*}
\frac{r^{\prime}}{r^{2}}=a(t)+b(t) r^{n-2} \tag{14}
\end{equation*}
$$

- Integrating equation (14) from $t_{0}$ to $t$ yields

$$
\begin{equation*}
r(t)=\frac{c}{1-c A(t)-c \int_{t_{0}}^{t} b(\tau) r^{n-2} d \tau} \tag{15}
\end{equation*}
$$

where $A(t)=\int_{t_{0}}^{t} a(\tau) d \tau$.

- Rewrite equation (15) as

$$
\begin{equation*}
r(t)=c\left(1+A(t)+r(t) \int_{t_{0}}^{t} b(\tau) r^{n-2} d \tau\right) \tag{16}
\end{equation*}
$$

## A Nonlinear Operator

- Let $\mathcal{C}[0,1]$ denote the Banach space of all continuous functions on the interval $[0,1]$ with the norm $\|f\|=\max _{0 \leq t \leq 1}|f(t)|$. We define the operator [10]:

$$
T_{c}: \mathcal{C}[0,1] \rightarrow \mathcal{C}[0,1],
$$

## A Nonlinear Operator

- Let $\mathcal{C}[0,1]$ denote the Banach space of all continuous functions on the interval $[0,1]$ with the norm $\|f\|=\max _{0 \leq t \leq 1}|f(t)|$. We define the operator [10]:

$$
\begin{gathered}
T_{c}: \mathcal{C}[0,1] \rightarrow \mathcal{C}[0,1] \\
T_{c}(f)(t) \stackrel{\text { def }}{=} \frac{c}{1-c A(t)-c \int_{0}^{t} b(\tau) f(\tau)^{n-2} d \tau}
\end{gathered}
$$

## A Nonlinear Operator

- Let $\mathcal{C}[0,1]$ denote the Banach space of all continuous functions on the interval $[0,1]$ with the norm $\|f\|=\max _{0 \leq t \leq 1}|f(t)|$. We define the operator [10]:

$$
\begin{gathered}
T_{c}: \mathcal{C}[0,1] \rightarrow \mathcal{C}[0,1] \\
T_{c}(f)(t) \stackrel{\text { def }}{=} \frac{c}{1-c A(t)-c \int_{0}^{t} b(\tau) f(\tau)^{n-2} d \tau}
\end{gathered}
$$

for given $a, b \in \mathcal{C}[0,1]$ and $c \in \mathbb{R}$. Obviously, $T_{c}$ is well defined on an arbitrary bounded set of $\mathcal{C}[0,1]$ if $c$ is suitably small. Let us first observe some useful properties of $T_{c}$.

## A Nonlinear Operator

- Let $\mathcal{C}[0,1]$ denote the Banach space of all continuous functions on the interval $[0,1]$ with the norm $\|f\|=\max _{0 \leq t \leq 1}|f(t)|$. We define the operator [10]:

$$
\begin{gathered}
T_{c}: \mathcal{C}[0,1] \rightarrow \mathcal{C}[0,1] \\
T_{c}(f)(t) \stackrel{\text { def }}{=} \frac{c}{1-c A(t)-c \int_{0}^{t} b(\tau) f(\tau)^{n-2} d \tau}
\end{gathered}
$$

for given $a, b \in \mathcal{C}[0,1]$ and $c \in \mathbb{R}$. Obviously, $T_{c}$ is well defined on an arbitrary bounded set of $\mathcal{C}[0,1]$ if $c$ is suitably small. Let us first observe some useful properties of $T_{c}$.

## A Nonlinear Operator

- Let $\mathcal{C}[0,1]$ denote the Banach space of all continuous functions on the interval $[0,1]$ with the norm $\|f\|=\max _{0 \leq t \leq 1}|f(t)|$. We define the operator [10]:

$$
\begin{gathered}
T_{c}: \mathcal{C}[0,1] \rightarrow \mathcal{C}[0,1] \\
T_{c}(f)(t) \stackrel{\text { def }}{=} \frac{c}{1-c A(t)-c \int_{0}^{t} b(\tau) f(\tau)^{n-2} d \tau}
\end{gathered}
$$

for given $a, b \in \mathcal{C}[0,1]$ and $c \in \mathbb{R}$. Obviously, $T_{c}$ is well defined on an arbitrary bounded set of $\mathcal{C}[0,1]$ if $c$ is suitably small. Let us first observe some useful properties of $T_{c}$.
[10] Z. Feng, Z. angew. Math. Phys. under review.

## Property of Our Operator

## Lemma (1)

For $f \in \mathcal{C}[0,1]$ and $c \in \mathbb{R}$ with $\|f\| \leq M$ and $|c|<c_{0} \stackrel{\text { def }}{=}\left(\|a\|+\|b\| M^{n-2}\right)^{-1}$, $T_{c}(f)$ is well defined and differentiable, and satisfies

$$
\frac{d}{d t} T_{c}(f)(t)=a(t)\left[T_{c}(f)(t)\right]^{2}+b(t)\left[T_{c}(f)(t)\right]^{2} f(t)^{n-2}
$$

## Property of Our Operator

## Lemma (1)

For $f \in \mathcal{C}[0,1]$ and $c \in \mathbb{R}$ with $\|f\| \leq M$ and $|c|<c_{0} \stackrel{\text { def }}{=}\left(\|a\|+\|b\| M^{n-2}\right)^{-1}$, $T_{c}(f)$ is well defined and differentiable, and satisfies

$$
\frac{d}{d t} T_{c}(f)(t)=a(t)\left[T_{c}(f)(t)\right]^{2}+b(t)\left[T_{c}(f)(t)\right]^{2} f(t)^{n-2}
$$

Furthermore, we have an identity

$$
\begin{array}{r}
T_{c}(f)(t)-T_{c}(g)(t)=T_{c}(f)(t) T_{c}(g)(t) \int_{0}^{t} b(\tau)\left(f(\tau)^{n-2}-g(\tau)^{n-2}\right) d \tau \\
0 \leq t \leq 1
\end{array}
$$

## Property of Our Operator

## Lemma (1)

For $f \in \mathcal{C}[0,1]$ and $c \in \mathbb{R}$ with $\|f\| \leq M$ and $|c|<c_{0} \stackrel{\text { def }}{=}\left(\|a\|+\|b\| M^{n-2}\right)^{-1}$, $T_{c}(f)$ is well defined and differentiable, and satisfies

$$
\frac{d}{d t} T_{c}(f)(t)=a(t)\left[T_{c}(f)(t)\right]^{2}+b(t)\left[T_{c}(f)(t)\right]^{2} f(t)^{n-2}
$$

Furthermore, we have an identity

$$
\begin{array}{r}
T_{c}(f)(t)-T_{c}(g)(t)=T_{c}(f)(t) T_{c}(g)(t) \int_{0}^{t} b(\tau)\left(f(\tau)^{n-2}-g(\tau)^{n-2}\right) d \tau \\
0 \leq t \leq 1
\end{array}
$$

for arbitrary $f, g \in \mathcal{C}[0,1]$ and $c \in \mathcal{R}$ with $\|f\|,\|g\| \leq M$ and $|c|<c_{0}$.

## Outline of the Proof

## Step 1: well-defined

$$
\begin{aligned}
& 1-c A(t)-c \int_{0}^{t} b(\tau) f(\tau)^{n-2} d \tau=0 \Rightarrow \\
& |c| \geq \frac{1}{|A(t)|+\int_{0}^{t}\left|b(\tau) f(\tau)^{n-2}\right| d \tau} \geq \frac{1}{\|a\|+\|b\| M^{n-2}}
\end{aligned}
$$

## Outline of the Proof

## Step 1: well-defined

$$
\begin{aligned}
& 1-c A(t)-c \int_{0}^{t} b(\tau) f(\tau)^{n-2} d \tau=0 \Rightarrow \\
& |c| \geq \frac{1}{|A(t)|+\int_{0}^{t}\left|b(\tau) f(\tau)^{n-2}\right| d \tau} \geq \frac{1}{\|a\|+\|b\| M^{n-2}}
\end{aligned}
$$

Step 2: A direct calculation gives

$$
\begin{aligned}
& \frac{d}{d t} T_{c}(f)(t)=\frac{-c\left[-c a(t)-c b(t) f(t)^{n-2}\right]}{\left(1-c A(t)-c \int_{0}^{t} b(\tau) f(\tau)^{n-2} d \tau\right)^{2}} \\
&= \frac{c^{2} a(t)}{\left(1-c A(t)-c \int_{0}^{t} b(\tau) f(\tau)^{n-2} d \tau\right)^{2}}+\frac{c^{2} b(t) f(t)^{n-2}}{\left(1-c A(t)-c \int_{0}^{t} b(\tau) f(\tau)^{n-2} d \tau\right)^{2}} \\
& T_{c}(f)(t)-T_{c}(g)(t)=\frac{c}{H(f)} \cdot \frac{c}{H(g)} \cdot \int_{0}^{t} b(\tau)\left(f(\tau)^{n-2}-g(\tau)^{n-2}\right) d \tau
\end{aligned}
$$

## Lemma 2

## Lemma (2)

Let $c_{1}=(\|a\|+\|b\|+1)^{-1}$. Then we have

$$
\left\|T_{c} f\right\| \leq 1 \quad \text { if }\|f\| \leq 1 \text { and }|c| \leq c_{1}
$$

## Lemma 2

## Lemma (2)

Let $c_{1}=(\|a\|+\|b\|+1)^{-1}$. Then we have

$$
\left\|T_{c} f\right\| \leq 1 \quad \text { if }\|f\| \leq 1 \text { and }|c| \leq c_{1}
$$

## Outline of the Proof.

If $|f| \mid \leq 1$ and $|c| \leq c_{1}$, then we have

$$
\begin{aligned}
\left\|T_{c} f\right\| & \leq \frac{|c|}{1-|c|\left(\|a\|+\|b\|\|f\|^{n-2}\right)} \\
& \leq \frac{|c|}{1-|c|(\|a\|+\|b\|)} \\
& \leq 1
\end{aligned}
$$

The conclusion follows.

## Lemma 3

## Lemma (3)

Let $c_{2}=(\sqrt{(n-2)\|b\|}+\|a\|+\|b\|+1)^{-1}$. If $|c| \leq c_{2}$, then $T_{c}$ is a contraction mapping on the close unit ball $\mathcal{B}_{1}=\{f \in \mathcal{C}[0,1] \mid\|f\| \leq 1\}$ of $\mathcal{C}[0,1]$.

## Outline of the Proof.

It follows from Lemmas 1 and 2 that

$$
\begin{aligned}
\left\|T_{c}(f)(t)-T_{c}(g)(t)\right\| \leq & \left\|T_{c}(f)\right\|\left\|T_{c}(g)\right\|\|b\|\left\|f^{n-2}-g^{n-2}\right\| \\
= & C\left\|(f-g)\left(f^{n-3}+f^{n-4} g+\cdots+f g^{n-4}+g^{n-3}\right)\right\| \\
& \leq(n-2) c\|f-g\|,
\end{aligned}
$$

where

$$
c \stackrel{\text { def }}{=}\left(\frac{|c|}{1-|c|(\|a\|+\|b\|)}\right)^{2}\|b\| .
$$

## Theorem 1

## Theorem (1)

For given $a, b \in \mathcal{C}[0,1]$ and $c \in \mathbb{R}$ with $|c| \leq(\sqrt{(n-2)\|b\|}+\|a\|+\|b\|+1)^{-1}$, the solution $r(t, c)$ of equation with $r(0, c)=c$ can be uniformly approximated by an iterated sequence $\left\{T_{c}^{n}(f)(t)\right\}:$

$$
\begin{equation*}
r(t, c)=\lim _{n \rightarrow \infty} T_{c}^{n}(f)(t), \quad 0 \leq t \leq 1 \tag{17}
\end{equation*}
$$

that is,

$$
\begin{equation*}
r(t, c)=\frac{c}{1-c A(t)-c^{n-1} \int_{0}^{t} \frac{b\left(t_{1}\right) d t_{1}}{b\left(t_{2}\right) d t_{2}}} \frac{1-c A\left(t_{1}\right)-c^{n-1} \int_{0}^{t_{1}} \frac{1-c A\left(t_{2}\right)-c^{n-1} \int_{0}^{t_{2}} \cdots}{1-1}}{1 .} \tag{18}
\end{equation*}
$$

for arbitrary $f \in \mathcal{C}[0,1]$ with $\|f\| \leq 1$. Furthermore, the following error estimate holds

$$
r(t, c)-T_{c}^{n}(f)(t)=\mathcal{O}\left(c^{2 n}\right)
$$

## Theorem 2: Case of $n=3$

- Denote

$$
M=\max _{t \in[0,1]}|a(t) \pm b(t)| .
$$

## Theorem 2: Case of $n=3$

- Denote

$$
M=\max _{t \in[0,1]}|a(t) \pm b(t)| .
$$

## Theorem (2)

Suppose $a, b \in \mathcal{C}[0,1]$ and $c \in \mathbb{R}$ with

$$
|c| \leq \max \left\{(\sqrt{\|b\|}+\|a\|+\|b\|+1)^{-1},(2 M)^{-1}\right\}
$$

Then, in formula (18), the following part is bounded

$$
\begin{aligned}
& \frac{b\left(t_{1}\right)}{1-c A\left(t_{1}\right)-c^{2} \int_{0}^{t_{1}} \frac{b\left(t_{2}\right) d t_{2}}{1-c A\left(t_{2}\right)-c^{2} \int_{0}^{t_{2} \ldots}}} \\
= & \frac{1}{c} \cdot b\left(t_{1}\right) \cdot \frac{c}{1-c A\left(t_{1}\right)-c^{2} \int_{0}^{t_{1}} b\left(t_{2}\right) \cdot \frac{c}{1-c A\left(t_{2}\right)-c^{2} \int_{0}^{t_{2} \ldots}} d t_{2}} .
\end{aligned}
$$

## 2D Korteweg-de Vries-Burgers Equation

- Consider the 2D Korteweg-de Vries-Burgers equation:

$$
\begin{equation*}
\left(U_{t}+\alpha U U_{x}+\beta U_{x x}+s U_{x x x}\right)_{x}+\gamma U_{y y}=0, \tag{19}
\end{equation*}
$$

## 2D Korteweg-de Vries-Burgers Equation

- Consider the 2D Korteweg-de Vries-Burgers equation:

$$
\begin{equation*}
\left(U_{t}+\alpha U U_{x}+\beta U_{x x}+s U_{x x x}\right)_{x}+\gamma U_{y y}=0, \tag{19}
\end{equation*}
$$

where $\alpha, \beta, s$, and $\gamma$ are constants and $\alpha \beta s \gamma \neq 0$.

## 2D Korteweg-de Vries-Burgers Equation

- Consider the 2D Korteweg-de Vries-Burgers equation:

$$
\begin{equation*}
\left(U_{t}+\alpha U U_{x}+\beta U_{x x}+s U_{x x x}\right)_{x}+\gamma U_{y y}=0, \tag{19}
\end{equation*}
$$

where $\alpha, \beta, s$, and $\gamma$ are constants and $\alpha \beta s \gamma \neq 0$.

- Assume that equation (19) has an exact solution in the form

$$
\begin{equation*}
U(x, y, t)=U(\xi), \quad \xi=h x+l y-w t . \tag{20}
\end{equation*}
$$

## 2D Korteweg-de Vries-Burgers Equation

- Consider the 2D Korteweg-de Vries-Burgers equation:

$$
\begin{equation*}
\left(U_{t}+\alpha U U_{x}+\beta U_{x x}+s U_{x x x}\right)_{x}+\gamma U_{y y}=0, \tag{19}
\end{equation*}
$$

where $\alpha, \beta, s$, and $\gamma$ are constants and $\alpha \beta s \gamma \neq 0$.

- Assume that equation (19) has an exact solution in the form

$$
\begin{equation*}
U(x, y, t)=U(\xi), \quad \xi=h x+l y-w t . \tag{20}
\end{equation*}
$$

- Substitution of (20) into equation (19) and performing integration twice yields

$$
\begin{gather*}
U^{\prime \prime}(\xi)+\lambda U^{\prime}(\xi)+a U^{2}(\xi)+b U(\xi)+d=0  \tag{21}\\
\text { where } v=U(\xi) \in\left[v_{0}, v_{1}\right], \lambda=\frac{\beta}{s h}, a=\frac{\alpha}{2 s h^{2}}, b=\frac{\gamma l^{2}-w h}{s h^{4}} \text { and } d=-\frac{C}{s h^{4}} .
\end{gather*}
$$

## Resultant Abel Equation

- Let $v=U(\xi)$ and $y=U^{\prime}(\xi)$. Equation (21) becomes

$$
\begin{equation*}
\frac{d y}{d v} y+\lambda y+a v^{2}+b v+d=0 . \tag{22}
\end{equation*}
$$

## Resultant Abel Equation

- Let $v=U(\xi)$ and $y=U^{\prime}(\xi)$. Equation (21) becomes

$$
\begin{equation*}
\frac{d y}{d v} y+\lambda y+a v^{2}+b v+d=0 . \tag{22}
\end{equation*}
$$

Using $z=\frac{1}{y}$ yields

$$
\begin{equation*}
\frac{d z}{d v}=\lambda z^{2}+\left(a v^{2}+b v+d\right) z^{3}, \quad z\left(v_{0}\right)=\frac{1}{U^{\prime}\left(\xi_{0}\right)}=c . \tag{23}
\end{equation*}
$$

## Resultant Abel Equation

- Let $v=U(\xi)$ and $y=U^{\prime}(\xi)$. Equation (21) becomes

$$
\begin{equation*}
\frac{d y}{d v} y+\lambda y+a v^{2}+b v+d=0 . \tag{22}
\end{equation*}
$$

Using $z=\frac{1}{y}$ yields

$$
\begin{equation*}
\frac{d z}{d v}=\lambda z^{2}+\left(a v^{2}+b v+d\right) z^{3}, \quad z\left(v_{0}\right)=\frac{1}{U^{\prime}\left(\xi_{0}\right)}=c . \tag{23}
\end{equation*}
$$

- Let $\eta=\frac{v-v_{0}}{v_{1}-v_{0}}$, then $\eta \in[0,1]$ and $v=v_{0}+\left(v_{1}-v_{0}\right) \eta$.


## Resultant Abel Equation

- Let $v=U(\xi)$ and $y=U^{\prime}(\xi)$. Equation (21) becomes

$$
\begin{equation*}
\frac{d y}{d v} y+\lambda y+a v^{2}+b v+d=0 . \tag{22}
\end{equation*}
$$

Using $z=\frac{1}{y}$ yields

$$
\begin{equation*}
\frac{d z}{d v}=\lambda z^{2}+\left(a v^{2}+b v+d\right) z^{3}, \quad z\left(v_{0}\right)=\frac{1}{U^{\prime}\left(\xi_{0}\right)}=c . \tag{23}
\end{equation*}
$$

- Let $\eta=\frac{v-v_{0}}{v_{1}-v_{0}}$, then $\eta \in[0,1]$ and $v=v_{0}+\left(v_{1}-v_{0}\right) \eta$. So equation (23) reduces to

$$
\begin{equation*}
r^{\prime}=h(\eta) r^{2}+k(\eta) r^{3}, \quad r(0)=c, \tag{24}
\end{equation*}
$$

## Resultant Abel Equation

- Let $v=U(\xi)$ and $y=U^{\prime}(\xi)$. Equation (21) becomes

$$
\begin{equation*}
\frac{d y}{d v} y+\lambda y+a v^{2}+b v+d=0 . \tag{22}
\end{equation*}
$$

Using $z=\frac{1}{y}$ yields

$$
\begin{equation*}
\frac{d z}{d v}=\lambda z^{2}+\left(a v^{2}+b v+d\right) z^{3}, \quad z\left(v_{0}\right)=\frac{1}{U^{\prime}\left(\xi_{0}\right)}=c . \tag{23}
\end{equation*}
$$

- Let $\eta=\frac{v-v_{0}}{v_{1}-v_{0}}$, then $\eta \in[0,1]$ and $v=v_{0}+\left(v_{1}-v_{0}\right) \eta$. So equation (23) reduces to

$$
\begin{equation*}
r^{\prime}=h(\eta) r^{2}+k(\eta) r^{3}, \quad r(0)=c \tag{24}
\end{equation*}
$$

where $h(\eta), k(\eta) \in \mathcal{C}[0,1]$, and

$$
\begin{gathered}
h(\eta)=\left(v_{1}-v_{0}\right) \lambda \\
k(\eta)=\left(v_{1}-v_{0}\right)\left(a v^{2}+b v+d\right)
\end{gathered}
$$

## Solution to Equation (24)

- By virtue of Theorem 1, if $|c| \leq(\sqrt{\|k\|}+\|h\|+\|k\|+1)^{-1}$, the solution to equation (24) is

$$
\begin{equation*}
r(\eta)=\lim _{n \rightarrow+\infty} T_{c}^{n}(w)(\eta) \tag{25}
\end{equation*}
$$

## Solution to Equation (24)

- By virtue of Theorem 1, if $|c| \leq(\sqrt{\|k\|}+\|h\|+\|k\|+1)^{-1}$, the solution to equation (24) is

$$
\begin{equation*}
r(\eta)=\lim _{n \rightarrow+\infty} T_{c}^{n}(w)(\eta) \tag{25}
\end{equation*}
$$

where $0 \leq \eta \leq 1$ for any $w \in \mathcal{C}[0,1]$ with $\|w\| \leq 1$, and

$$
T_{c}(w)=\frac{c}{1-c H(\eta)-c \int_{0}^{\eta} k(x) w(x)^{n-2} d x}
$$

## Solution to Equation (24)

- By virtue of Theorem 1, if $|c| \leq(\sqrt{\|k\|}+\|h\|+\|k\|+1)^{-1}$, the solution to equation (24) is

$$
\begin{equation*}
r(\eta)=\lim _{n \rightarrow+\infty} T_{c}^{n}(w)(\eta) \tag{25}
\end{equation*}
$$

where $0 \leq \eta \leq 1$ for any $w \in \mathcal{C}[0,1]$ with $\|w\| \leq 1$, and

$$
T_{c}(w)=\frac{c}{1-c H(\eta)-c \int_{0}^{\eta} k(x) w(x)^{n-2} d x}
$$

where

$$
\begin{aligned}
H(\eta) & =\int_{0}^{\eta} h(x) d x=\int_{0}^{\eta}\left(v_{1}-v_{0}\right) \lambda d x=\left(v_{1}-v_{0}\right) \lambda \eta \\
k(x) & =\left(v_{1}-v_{0}\right)\left(a\left(v_{0}+\left(v_{1}-v_{0}\right) x\right)^{2}+b\left(v_{0}+\left(v_{1}-v_{0}\right) x\right)+d\right)
\end{aligned}
$$

## Approximate Solution to 2D-KdV-Burgers Equation

- Recall that $r=\frac{1}{y}, y=U^{\prime}(\xi), \eta=\frac{v-v_{0}}{v_{1}-v_{0}}$ and $v=U(\xi)$. When conditions of Theorem 1 are fulfilled, we have

$$
\begin{equation*}
\frac{1}{U^{\prime}(\xi)}=\frac{c}{1-c A(\xi)-c^{2} \int_{0}^{\xi} \frac{b\left(t_{1}\right) d t_{1}}{1-c A\left(t_{1}\right)-c^{2} \int_{0}^{t_{1}} \frac{b\left(t_{2}\right) d t_{2}}{1-c A\left(t_{2}\right)-c^{2} \int_{0}^{t_{2}} \cdots}}} . \tag{26}
\end{equation*}
$$

## Approximate Solution to 2D-KdV-Burgers Equation

- Recall that $r=\frac{1}{y}, y=U^{\prime}(\xi), \eta=\frac{v-v_{0}}{v_{1}-v_{0}}$ and $v=U(\xi)$. When conditions of Theorem 1 are fulfilled, we have

$$
\begin{equation*}
\frac{1}{U^{\prime}(\xi)}=\frac{c}{1-c A(\xi)-c^{2} \int_{0}^{\xi} \frac{b\left(t_{1}\right) d t_{1}}{1-c A\left(t_{1}\right)-c^{2} \int_{0}^{t_{1}} \frac{b\left(t_{2}\right) d t_{2}}{1-c A\left(t_{2}\right)-c^{2} \int_{0}^{t_{2}} \cdots}}} . \tag{26}
\end{equation*}
$$

- When $c$ is small, according to Theorem 2, the coefficient of $c^{2}$ is bounded. So we can drop the term containing $c^{2}$ and get

$$
\begin{aligned}
U^{\prime}(\xi) & \approx \frac{1-c\left(v_{1}-v_{0}\right) \lambda \eta}{c} \\
& =\frac{1-c \lambda\left(U(\xi)-v_{0}\right)}{c}
\end{aligned}
$$

## Approximate Solution to 2D-KdV-Burgers Equation

- Recall that $r=\frac{1}{y}, y=U^{\prime}(\xi), \eta=\frac{v-v_{0}}{v_{1}-v_{0}}$ and $v=U(\xi)$. When conditions of Theorem 1 are fulfilled, we have

$$
\begin{equation*}
\frac{1}{U^{\prime}(\xi)}=\frac{c}{1-c A(\xi)-c^{2} \int_{0}^{\xi} \frac{b\left(t_{1}\right) d t_{1}}{1-c A\left(t_{1}\right)-c^{2} \int_{0}^{t_{1}} \frac{b\left(t_{2}\right) d t_{2}}{1-c A\left(t_{2}\right)-c^{2} \int_{0}^{t_{2} \cdots}}}} . \tag{26}
\end{equation*}
$$

- When $c$ is small, according to Theorem 2, the coefficient of $c^{2}$ is bounded. So we can drop the term containing $c^{2}$ and get

$$
\begin{aligned}
U^{\prime}(\xi) & \approx \frac{1-c\left(v_{1}-v_{0}\right) \lambda \eta}{c} \\
& =\frac{1-c \lambda\left(U(\xi)-v_{0}\right)}{c}
\end{aligned}
$$

That is,

$$
\begin{equation*}
U^{\prime}(\xi)+\lambda U(\xi)=\frac{1}{c}+\lambda v_{0} . \tag{27}
\end{equation*}
$$

## Approximate Solution to 2D KdV-Burgers Equation

- Solving equation (27) gives

$$
U(x, y, t)=\frac{\frac{1}{c}+\lambda v_{0}}{\lambda}+c e^{-\lambda \xi}, \quad \xi=h x+l y-w t
$$

where $\lambda=\frac{\beta}{s h}$.

## Approximate Solution to 2D KdV-Burgers Equation

- Solving equation (27) gives

$$
U(x, y, t)=\frac{\frac{1}{c}+\lambda v_{0}}{\lambda}+c e^{-\lambda \xi}, \quad \xi=h x+l y-w t
$$

where $\lambda=\frac{\beta}{s h}$.

- If we take $v_{0}=\frac{b}{2 a}$ and choose $c=\frac{-2 a}{\lambda \sqrt{b^{2}-4 a d}}$ sufficiently small, when $\lambda \xi \rightarrow+\infty$, we obtain

$$
\begin{equation*}
U(x, y, t) \sim \frac{b^{2}-4 a d}{-2 a}+\frac{b}{2 a} . \tag{28}
\end{equation*}
$$

## Approximate Solution to 2D KdV-Burgers Equation

- Solving equation (27) gives

$$
U(x, y, t)=\frac{\frac{1}{c}+\lambda v_{0}}{\lambda}+c e^{-\lambda \xi}, \quad \xi=h x+l y-w t
$$

where $\lambda=\frac{\beta}{s h}$.

- If we take $v_{0}=\frac{b}{2 a}$ and choose $c=\frac{-2 a}{\lambda \sqrt{b^{2}-4 a d}}$ sufficiently small, when $\lambda \xi \rightarrow+\infty$, we obtain

$$
\begin{equation*}
U(x, y, t) \sim \frac{b^{2}-4 a d}{-2 a}+\frac{b}{2 a} . \tag{28}
\end{equation*}
$$

- It is remarkable that the approximate solution (28) is in agreement with main results described in $[7,8]$ by the Hardy's theory and the theory of Lie symmetry.


## Approximate Solution to 2D KdV-Burgers Equation

- Solving equation (27) gives

$$
U(x, y, t)=\frac{\frac{1}{c}+\lambda v_{0}}{\lambda}+c e^{-\lambda \xi}, \quad \xi=h x+l y-w t
$$

where $\lambda=\frac{\beta}{s h}$.

- If we take $v_{0}=\frac{b}{2 a}$ and choose $c=\frac{-2 a}{\lambda \sqrt{b^{2}-4 a d}}$ sufficiently small, when $\lambda \xi \rightarrow+\infty$, we obtain

$$
\begin{equation*}
U(x, y, t) \sim \frac{b^{2}-4 a d}{-2 a}+\frac{b}{2 a} . \tag{28}
\end{equation*}
$$

- It is remarkable that the approximate solution (28) is in agreement with main results described in $[7,8]$ by the Hardy's theory and the theory of Lie symmetry.


## Approximate Solution to 2D KdV-Burgers Equation

- Solving equation (27) gives

$$
U(x, y, t)=\frac{\frac{1}{c}+\lambda v_{0}}{\lambda}+c e^{-\lambda \xi}, \quad \xi=h x+l y-w t
$$

where $\lambda=\frac{\beta}{s h}$.

- If we take $v_{0}=\frac{b}{2 a}$ and choose $c=\frac{-2 a}{\lambda \sqrt{b^{2}-4 a d}}$ sufficiently small, when $\lambda \xi \rightarrow+\infty$, we obtain

$$
\begin{equation*}
U(x, y, t) \sim \frac{b^{2}-4 a d}{-2 a}+\frac{b}{2 a} \tag{28}
\end{equation*}
$$

- It is remarkable that the approximate solution (28) is in agreement with main results described in $[7,8]$ by the Hardy's theory and the theory of Lie symmetry.
[7] Z. Feng, J. Phys. A (Math. Gen.) 36 (2003), 8817-8827.


## Approximate Solution to 2D KdV-Burgers Equation

- Solving equation (27) gives

$$
U(x, y, t)=\frac{\frac{1}{c}+\lambda v_{0}}{\lambda}+c e^{-\lambda \xi}, \quad \xi=h x+l y-w t
$$

where $\lambda=\frac{\beta}{s h}$.

- If we take $v_{0}=\frac{b}{2 a}$ and choose $c=\frac{-2 a}{\lambda \sqrt{b^{2}-4 a d}}$ sufficiently small, when $\lambda \xi \rightarrow+\infty$, we obtain

$$
\begin{equation*}
U(x, y, t) \sim \frac{b^{2}-4 a d}{-2 a}+\frac{b}{2 a} . \tag{28}
\end{equation*}
$$

- It is remarkable that the approximate solution (28) is in agreement with main results described in $[7,8]$ by the Hardy's theory and the theory of Lie symmetry.
[7] Z. Feng, J. Phys. A (Math. Gen.) 36 (2003), 8817-8827.
[8] Z. Feng, Nonlinearity, 20 (2007), 343-356.


## Boundedness of Solutions

- Note that equation (26) can be rewritten as

$$
\begin{equation*}
\frac{1}{U^{\prime}(\xi)}=\frac{c}{1-c A(\xi)-c^{2} \Phi(\xi)} \tag{29}
\end{equation*}
$$

where $L \leq \Phi(\xi) \leq R$.

## Boundedness of Solutions

- Note that equation (26) can be rewritten as

$$
\begin{equation*}
\frac{1}{U^{\prime}(\xi)}=\frac{c}{1-c A(\xi)-c^{2} \Phi(\xi)} \tag{29}
\end{equation*}
$$

where $L \leq \Phi(\xi) \leq R$.

- When $\Phi$ is a quadratic or cubic function or special function of $U(\xi)$, we can analyze equation (29) qualitatively and numerically with classifications. For instance, if $\Phi$ is quadratic, we take $v_{0}=\frac{b}{2 a}$ and choose $c=\frac{-2 a}{\lambda \sqrt{b^{2}-4 a d}}$ sufficiently small, we can obtain the solution of the type

$$
u(x, y, t)=\frac{3 \beta^{2}+\gamma+c}{25 \alpha s} \operatorname{sech}^{2} \xi-\frac{6 \beta^{2}+\gamma+c}{25 \alpha s} \tanh \xi \pm \frac{6 \beta^{2}}{25 \alpha s}+C_{0}
$$

## Boundedness of Solutions

- Note that equation (26) can be rewritten as

$$
\begin{equation*}
\frac{1}{U^{\prime}(\xi)}=\frac{c}{1-c A(\xi)-c^{2} \Phi(\xi)} \tag{29}
\end{equation*}
$$

where $L \leq \Phi(\xi) \leq R$.

- When $\Phi$ is a quadratic or cubic function or special function of $U(\xi)$, we can analyze equation (29) qualitatively and numerically with classifications. For instance, if $\Phi$ is quadratic, we take $v_{0}=\frac{b}{2 a}$ and choose $c=\frac{-2 a}{\lambda \sqrt{b^{2}-4 a d}}$ sufficiently small, we can obtain the solution of the type
$u(x, y, t)=\frac{3 \beta^{2}+\gamma+c}{25 \alpha s} \operatorname{sech}^{2} \xi-\frac{6 \beta^{2}+\gamma+c}{25 \alpha s} \tanh \xi \pm \frac{6 \beta^{2}}{25 \alpha s}+C_{0}$.
- When $\Phi$ is a function with the lower and upper bounds, we can also find bounds of solutions of equation (29) by the comparison principle, which match well with the phase analysis described in [7].


## Summary

- In this talk, we provided a connection between the Abel equation of the first kind, an ordinary differential equation that is cubic in the unknown function, and the Korteweg-de Vries-Burgers equation, a partial differential equation that describes the propagation of waves on liquid-filled elastic tubes. We presented an integral form of the Abel equation with the initial condition.


## Summary

- In this talk, we provided a connection between the Abel equation of the first kind, an ordinary differential equation that is cubic in the unknown function, and the Korteweg-de Vries-Burgers equation, a partial differential equation that describes the propagation of waves on liquid-filled elastic tubes. We presented an integral form of the Abel equation with the initial condition.
- By virtue of the integral form and the Banach Contraction Mapping Principle we derived the asymptotic expansion of bounded solutions in the Banach space, and used the asymptotic formula to construct approximate solutions to the Korteweg-de Vries-Burgers equation.


## Summary

- In this talk, we provided a connection between the Abel equation of the first kind, an ordinary differential equation that is cubic in the unknown function, and the Korteweg-de Vries-Burgers equation, a partial differential equation that describes the propagation of waves on liquid-filled elastic tubes. We presented an integral form of the Abel equation with the initial condition.
- By virtue of the integral form and the Banach Contraction Mapping Principle we derived the asymptotic expansion of bounded solutions in the Banach space, and used the asymptotic formula to construct approximate solutions to the Korteweg-de Vries-Burgers equation.
- As an example, we presented the asymptotic behavior of traveling wave solution for a 2D KdV-Burgers equation which agrees well with existing results in the literature.


## Summary

- In this talk, we provided a connection between the Abel equation of the first kind, an ordinary differential equation that is cubic in the unknown function, and the Korteweg-de Vries-Burgers equation, a partial differential equation that describes the propagation of waves on liquid-filled elastic tubes. We presented an integral form of the Abel equation with the initial condition.
- By virtue of the integral form and the Banach Contraction Mapping Principle we derived the asymptotic expansion of bounded solutions in the Banach space, and used the asymptotic formula to construct approximate solutions to the Korteweg-de Vries-Burgers equation.
- As an example, we presented the asymptotic behavior of traveling wave solution for a 2D KdV-Burgers equation which agrees well with existing results in the literature.
- Under certain conditions, we can also study bounds of traveling wave solutions of KdV-Burgers type equations by the comparison principle.


## Acknowledgement

- I would like to thank


## Acknowledgement

- I would like to thank Xiaoqian Gong


## Acknowledgement

- I would like to thank Xiaoqian Gong for discussions and help on computations.


## Acknowledgement

- I would like to thank Xiaoqian Gong for discussions and help on computations.
- Thank you.

