

Representation Theory as Gauge Theory

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Themes

I. Harmonic analysis as the exploitation of symmetry¹

II. Commutative algebra signals geometry

III. Topology provides commutativity

IV. Gauge theory bridges topology and representation theory

¹Mackey, Bull. AMS 1980

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Outline

Representation theory

Quantum Field Theory

Gauge Theory

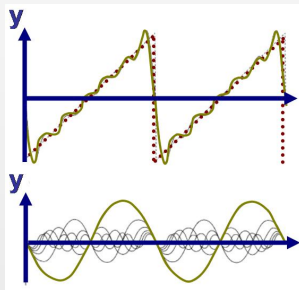
Fourier Series

$G = U(1)$ acts on S^1 , hence on $L^2(S^1)$:

Fourier series

$$L^2(S^1) \simeq \widehat{\bigoplus_{n \in \mathbb{Z}} \mathbb{C} e^{2\pi i n \theta}}$$

joint eigenspaces of rotation operators



The dual

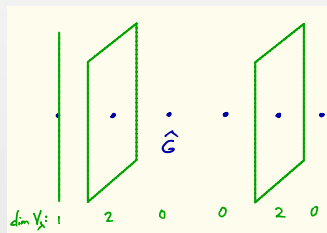
Basic object in representation theory:
describe the **dual** of a group G :

$$\widehat{G} = \{ \text{irreducible (unitary) representations of } G \}$$

e.g. $\widehat{U(1)} \simeq \mathbb{Z}$

$G \curvearrowright V$ representation \implies

spectral decomposition of V as family of vector spaces over dual

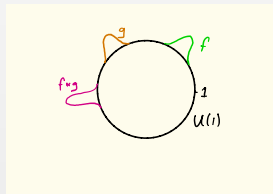


Group algebras

How to access \widehat{G} ?

Linear combinations of group elements form ring $\mathbb{C}G$
acting on any representation V

e.g., $C(U(1)) \curvearrowright L^2(S^1)$ (convolution)



Pontrjagin duality

G abelian

$\Rightarrow \mathbb{C}G$ is commutative (ring or C^* algebra)

$\Rightarrow \hat{G}$ underlies space: its spectrum

II. Commutative algebra signals geometry.

Spectral decomposition: $\mathbb{C}G$ -modules give
sheaves / vector bundles / projection valued measures
on \hat{G}

Fourier transform

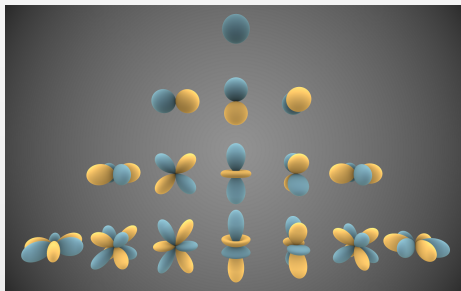
G	\hat{G}
Representation theory	Geometry
$U(1)$ (compact)	\mathbb{Z} (discrete)
\mathbb{R}	\mathbb{R}
Irreducible reps $\mathbb{C}e^{2\pi ixt}$	Points $t \in \mathbb{Z}$ or \mathbb{R}
$L^2(G)$ basis of irreducibles	$L^2(\hat{G})$ (Fourier transform)
Translation (group) Differentiation (enveloping algebra) Convolution (group algebra)	Multiplication (simultaneous diagonalization)
Representation	Family of vector spaces (spectral decomposition)

Example: Spherical Harmonics

$G = SO(3)$ acts on S^2 , hence on $L^2(S^2)$

$$L^2(S^2) \simeq \widehat{\bigoplus_{l \in \mathbb{Z}_+} V_l}$$

irreducible representations of $SO(3)$ (of $\dim 2l + 1$)



Example: spherical harmonics

$SO(3) \curvearrowright L^2(S^2)$ commutes with spherical Laplacian Δ :

Irreducible summands are precisely Δ -eigenspaces

Δ comes from quadratic **Casimir** $C = \frac{1}{2}(\mathbf{i}^2 + \mathbf{j}^2 + \mathbf{k}^2)$
in center of enveloping algebra of $SO(3)$, gives function on unitary
dual

$$\widehat{SO(3)} \simeq \mathbb{Z}_+, \quad C|_{V_l} = l(l+1)$$

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Nonabelian Duality

How to access \widehat{G} for G nonabelian?

seek commutative algebra.

Center of $\mathbb{C}G$: commute with all $g \in G$

\Rightarrow operators in G averaged over conjugacy classes:

class functions $\mathbb{C}G/G$

Schur's lemma:

Center acts by scalar on any irreducible representation:

give functions on \widehat{G}

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Dual and center

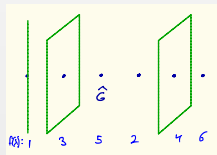
Conversely:

Functions on \widehat{G} give symmetries of arbitrary representations of G

– act as scalars on fibers of spectral decomposition

\rightsquigarrow operators commuting with G (and G -maps):

Bernstein center

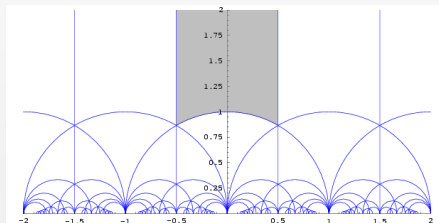


Automorphic Forms

Triumph of **exploitation of symmetry (I)**: the Langlands program.

Problem: spectral decomposition of $L^2(X)$ for X **moduli space of elliptic curves**

$$SL_2\mathbb{Z} \backslash \{\text{upper half plane}\} \simeq SL_2\mathbb{Z} \backslash SL_2\mathbb{R} / SO_2$$



more generally **arithmetic locally symmetric spaces**

$$X = \Gamma \backslash G_{\mathbb{R}} / K$$

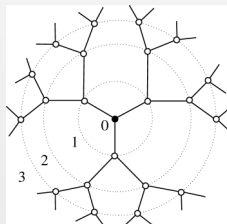
$G_{\mathbb{R}}$ real reductive group, $K \subset G_{\mathbb{R}}$ compact, $\Gamma \subset G_{\mathbb{R}}$ arithmetic
 (\leftrightarrow a number field F , and “level structure” at some primes)

Hecke operators

What acts?

- Laplacian Δ_X (and higher Casimirs): center of $U\mathfrak{g}$
- For every prime p , algebra of **Hecke operators**:
“higher discrete Laplacians”

e.g. average functions over elliptic curves “differing only at p ”
(by p -isogenies)



Holy grail of representation theory

All Hecke operators at almost all p commute



Langlands program:

1. Study spectral decomposition of $L^2(X)$ under Hecke operators
2. Identify joint spectrum with a space of Galois representations
3. Access mysteries of the universe

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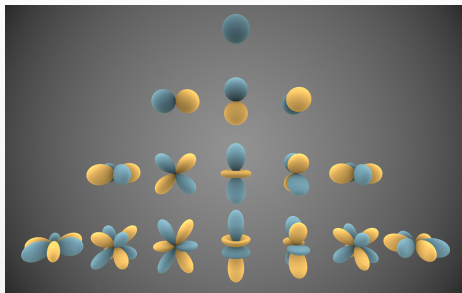
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Quantum Field Theory

Gauge Theory

The Gruppenpest



$L^2(X)$: Hilbert space of a quantum particle on X

$\frac{1}{2}\Delta_X$: Hamiltonian

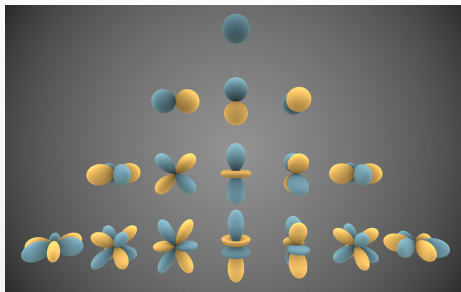
$G \curvearrowright X \implies G \curvearrowright L^2(X)$:

Symmetries of quantum system.

Exploitation of symmetry (I): find spectrum!

\rightsquigarrow The plague of groups.

The Gruppenpest



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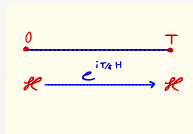
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Quantum mechanics: States

Quantum mechanics attaches

- Hilbert space \mathcal{H}



- Time evolution: group generated by Hamiltonian H

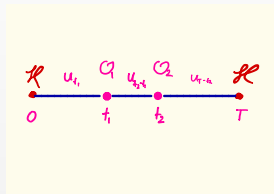
$$U_T = e^{iTH/\hbar} : \mathcal{H} \rightarrow \mathcal{H}$$



Quantum Mechanics: Operators

- Observables: (self-adjoint) operators on \mathcal{H} .

Measurement \leftrightarrow diagonalization



Noncommutative algebra (order in time matters!)

\rightsquigarrow can't diagonalize both position (x) and momentum ($-i\hbar \frac{\partial}{\partial x}$):

Heisenberg uncertainty

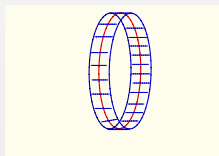


Sketch of Quantum Field Theory, after Segal



An n -dimensional² quantum field theory \mathcal{Z} gives rise to:

- “States”: A vector space $\mathcal{Z}(M)$ assigned to³ $(n-1)$ -dimensional Riemannian manifold M



Disjoint unions \mapsto tensor products

²(Euclidean)

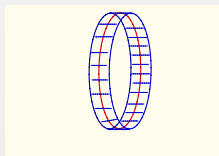
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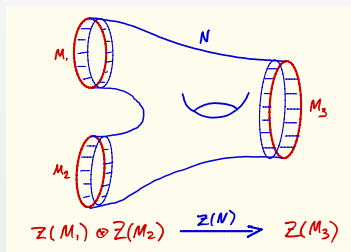
²(Euclidean)

³(n -dimensional collars around)

Cobordisms

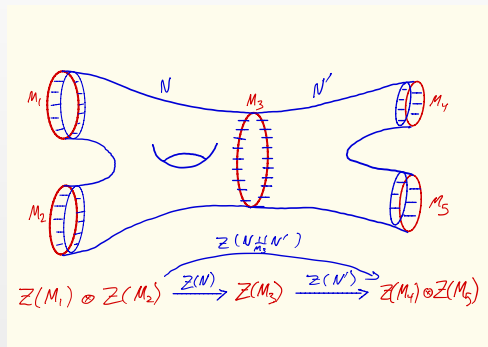
- “Time evolution”: An n -dimensional cobordism $N : M_{in} \rightarrow M_{out}$ defines a linear map

$$\mathcal{Z}(N) : \mathcal{Z}(M_{in}) \rightarrow \mathcal{Z}(M_{out}),$$



Composition

- **Time evolution:** Maps compose under gluing.



Topological Field Theory

Radical simplification: make \mathcal{Z} depend only on
topology of spacetime M

Mechanism: **Supersymmetry / Hodge theory**

e.g. in quantum mechanics ($n = 1$):
pass from functions $L^2(X)$ to forms $\Omega^\bullet(X)$,

$$\Delta = dd^* + d^*d$$

$\Rightarrow H = 0$ on cohomology:

Time evolution disappears, only space of states left:

Topological Quantum Mechanics \leftrightarrow a vector space

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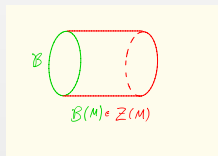
Extended locality

States $\mathcal{Z}(M)$ in TFT attached to compact $(n-1)$ -manifolds - complicated!

Avatars of states attached to small patches of space:

local boundary conditions

$(n-1)$ -dimensional field theories \mathcal{B}
which output states on compact M



Example: σ -models

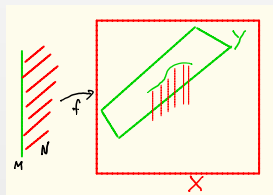
σ -model: Quantum theory of maps of space-time to target X

($n = 1$: quantum particle on X)

Local boundary conditions ($n = 2$):

quantum mechanical systems “over X ”, e.g.:

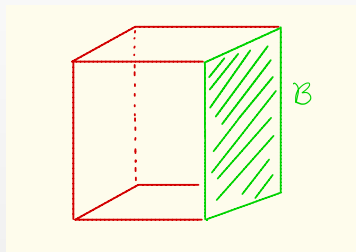
- quantum particle on a space $Y \longrightarrow X$
- family \mathcal{H} of vector spaces over X



Produce cohomology classes on X ($\sim \mathcal{Z}(S^1)$)
(fundamental class / Chern character)

Algebra of boundary conditions

What do boundary conditions form?

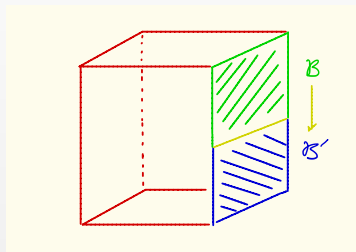


Boundary conditions can have interfaces (**morphisms**), and interfaces between interfaces (**2-morphisms**), and ... :

(n-1)-category

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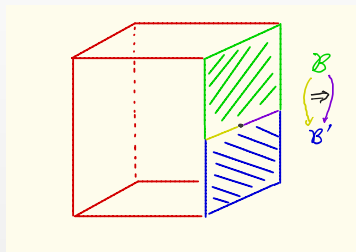


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How do you construct a TFT?

Topological QM determined by its states (vector space \mathcal{H})

Topological QFT determined by its local boundary conditions:

Theorem (Lurie)

Baez-Dolan Cobordism Hypothesis: An n -dimensional topological field theory \mathcal{Z} is uniquely determined by its higher category of boundary conditions.

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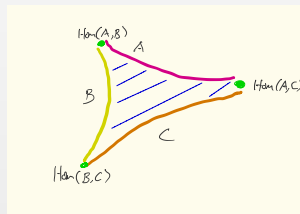
Cobordism Hypothesis

Cobordism Hypothesis:

- input (suitably finite) higher category
- output topological field theory

e.g.,

Categories \leftrightarrow 2-dimensional TFTs



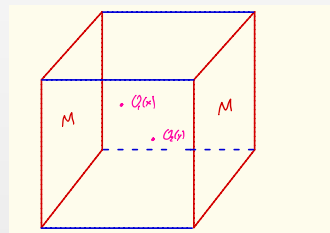
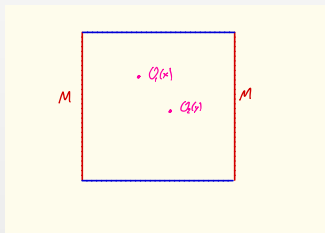
deep geometric perspective on categories

Observables



Key structure: **local operators**.

Measurements at points of space-time



Operators in TFT \leftrightarrow states on a small linking sphere, $\mathcal{Z}(S^{n-1})$.

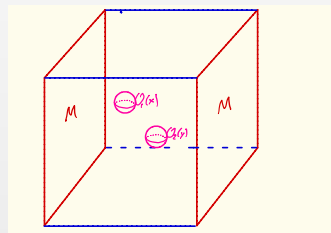
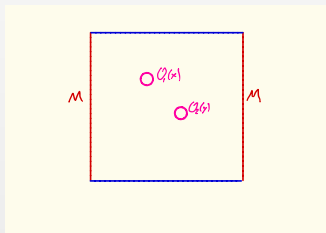
Act on states, and on boundary conditions

Observables



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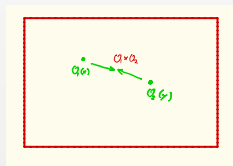
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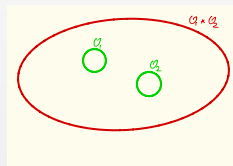
Operator product expansion

Topological pictures of composition of local operators:

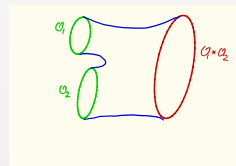
$$\mathcal{Z}(S^{n-1}) \otimes \mathcal{Z}(S^{n-1}) \longrightarrow \mathcal{Z}(S^{n-1})$$



OPE



Little Discs



Pair of Pants

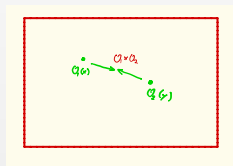
For $n > 1$, can move operators around each other:

local operators in TFT commute!

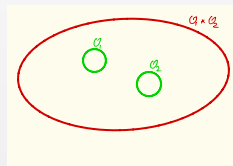
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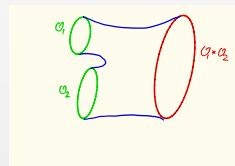
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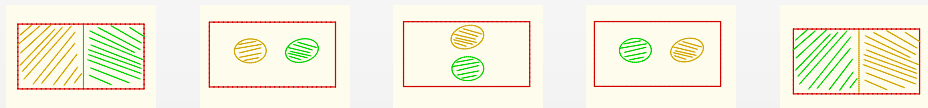
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Local operators commute

Commutativity of local operators:
same as commutativity of π_n ($n \geq 2$).

Picture of based maps to a pointed target $(X, *)$:



more room to move around!

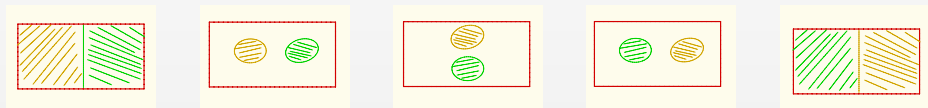
Noncommutativity of quantum mechanics \rightsquigarrow commutativity in TFT

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III. Topology provides commutativity

Fourier Transform for TFT

II. Commutative algebra signals geometry

Local operators in $\mathcal{Z} \implies$ moduli space of vacua $\mathfrak{M}_{\mathcal{Z}}$

\mathcal{Z}	$\mathfrak{M}_{\mathcal{Z}}$
TFT	Geometry σ -model low-energy effective theory
Local operators	Functions vacuum expectation values
OPE	Multiplication
Boundary conditions (action of local operators)	Families over $\mathfrak{M}_{\mathcal{Z}}$ (sheaves)
2d TFT	$\mathfrak{M}_{\mathcal{Z}}$ discrete

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Gauge Theory

Gauge theory

G compact Lie group

Gauge theory: Quantum theory of G -bundles with connection
(up to gauge transformations) on space-time

Tremendously influential in low dimensional topology:

invariants of manifolds M from cohomology of moduli spaces of G -bundles on M .

Revenge of the Gruppenpest

Local point of view:

Locally connections are trivial, but carry point-wise G symmetry

\implies boundary conditions: $(n-1)$ -dim QFTs with G symmetry:
representations of G on field theories.

Cobordism hypothesis: this determines the topological theory.

Quantum mechanics utilizes representation theory

Quantum field theory can **encode** representation theory

IV. Gauge theory bridges topology and representation theory

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IV. Gauge theory bridges topology and representation theory

2d Yang-Mills theory

2d Yang-Mills theory:

“Counts” G -bundles with connection on surfaces
(literally for G finite)

Boundary conditions: quantum mechanics with G -symmetry:

e.g. particle on a G -space X : $G \curvearrowright (L^2(X), \Delta_X)$:

Topological Yang-Mills: can define theory via

boundary conditions \leftrightarrow representations of G

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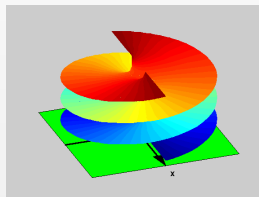
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Topological Yang-Mills: can **define** theory via

boundary conditions \leftrightarrow representations of G

Local operators

Local operators: functions on G -connections on S^1 ,
i.e. functions of **monodromy**, conjugacy class $[g] \in G/G$

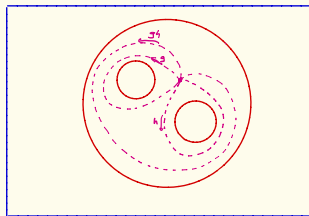


Local operators are class functions

$$\mathcal{Z}(S^1) = \mathbb{C}G/G$$

Local operators are Bernstein center

OPE \leftrightarrow convolution of class functions:



Local operators

\leftrightarrow **center** of group algebra $\mathbb{C}G$

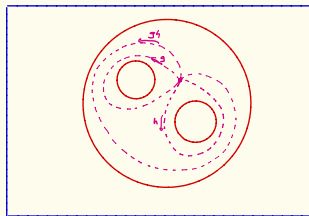
\leftrightarrow **Bernstein center** of boundary conditions

\leftrightarrow functions on dual \hat{G}

The dual \hat{G} is the moduli space of vacua $\mathfrak{M}_{\mathcal{Z}}$

Local operators are Bernstein center

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Local operators

\leftrightarrow **center** of group algebra $\mathbb{C}G$

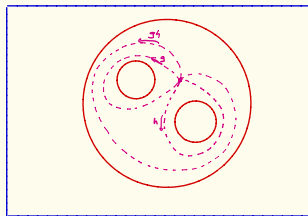
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Local operators

\leftrightarrow **center** of group algebra $\mathbb{C}G$

\leftrightarrow **Bernstein center** of boundary conditions

\leftrightarrow functions on dual \widehat{G}

The dual \widehat{G} is the moduli space of vacua $\mathfrak{M}_{\mathbb{Z}}$

Propaganda for three dimensions

3d gauge theories are precisely
modern geometric setting for representation theory

- **Seiberg-Witten Geometry:** Moduli spaces $\mathfrak{M}_{\mathbb{Z}}$ symplectic, carry canonical integrable systems, quantizations
(cf. Braverman's talk Monday)
- Teleman: powerful exploitation of symmetry in symplectic geometry (Gromov-Witten theory)
(talk Friday)
- BZ-Gunningham-Nadler: new commutative symmetries and spectral decomposition in classical representation theory
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Geometric Langlands Program

G complex reductive ($GL_n\mathbb{C}$, $SL_n\mathbb{C}$, $SO_n\mathbb{C}$, $Sp_n\mathbb{C}$,...)

Langlands	Geometric Langlands
Number field	Riemann surface C
Locally symmetric space $X_\Gamma = \Gamma \backslash G_\mathbb{R} / K$	Moduli of G -bundles on C $Bun_G(C)$
Hilbert space $L^2(X_\Gamma)$	"Orbifold" Cohomology of $Bun_G(C)$
Hecke operator at prime p	Hecke operator at point $x \in C$
modify elliptic curve by p -isogeny	modify G -bundle at x

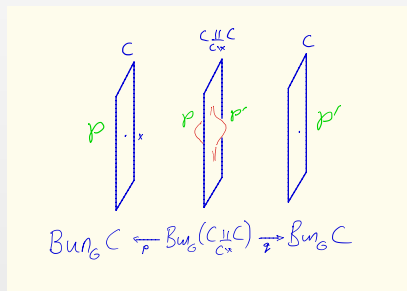
Modifying bundles

Hecke operators:

Convolution operators on $Bun_G(C)$:

take weighted average of a cohomology class
over bundles differing only by **modification at $x \in C$**

($\leftrightarrow \mathcal{P}, \mathcal{P}'$ glue to bundle on C with c doubled)



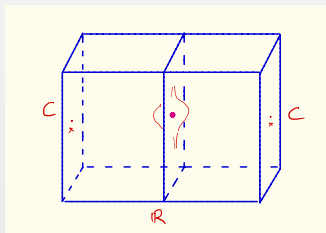
Hecke modifications as magnetic monopoles

Kapustin-Witten: interpretation via (maximally supersymmetric) gauge theory.

Hecke modification is creation of magnetic monopole in 3d theory on $C \times \mathbb{R}$

- prescribed singularity in gauge fields at point of monopole
- i.e., 't Hooft operator in gauge theory.

Acts as operator on states on C .



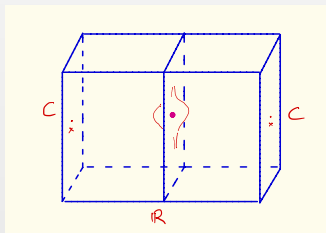
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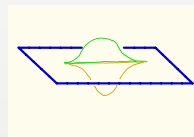
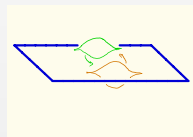
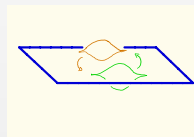
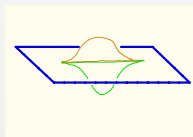
Acts as operator on states on C .



Why do Hecke operators commute?

Inspired by operator products in QFT, Beilinson-Drinfeld discovered source of commutativity:

to compose Hecke modifications, we can stack or slide them around the surface:



Commutativity of Hecke operators is that of
local operators and of homotopy groups (III)

Langlands Duality

Commutativity of Hecke operators signals geometry (II)

\implies calculate spectrum / dual / moduli space:

Geometric Satake Theorem⁴:

Hecke operators \leftrightarrow representations of
the Langlands dual group G^\vee

$$(GL_n \leftrightarrow GL_n, SL_n \leftrightarrow PGL_n, SO_{2n+1} \leftrightarrow Sp_{2n}, \dots)$$

\implies spectrum is **character variety**

$$Loc_{G^\vee}(C) = \{\pi_1(C) \rightarrow G^\vee\}$$

(geometric analog of Galois representations)

⁴Lusztig, Ginzburg, Mirkovic-Vilonen

Duality in physics

Kapustin-Witten: realize as **electric-magnetic duality**:
nonabelian, SUSY version⁵ of symmetry of Maxwell's equations
under exchange of electric and magnetic fields

G	G^\vee
Magnetic Automorphic	Electric Spectral
$Bun_G(C)$	G^\vee Character variety $\{\pi_1(C) \rightarrow G^\vee\}$
Creation of monopole ('t Hooft operator) (Hecke operator)	Measurement of field (Wilson operator) (measure holonomy)
Representation theory	Geometry

⁵Montonen-Olive S-duality

Factorization algebras

Beilinson-Drinfeld:

Factorization algebra:

algebraic structure encoded in collisions of points

↪ geometric theory of operators in

1. conformal field theory (vertex algebras),
2. topological field theory⁶, and
3. quantum field theory in general⁷

⁶Lurie, Francis-Ayala

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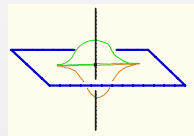
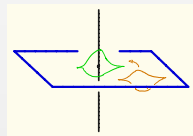
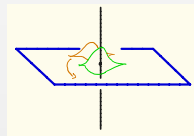
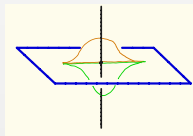
⁶Lurie, Francis-Ayala

⁷Costello-Gwilliam

Bernstein center

Ramification / level structure: allow singularities at points of C
(physically: solenoids)⁸

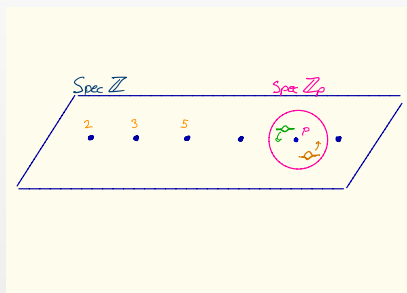
Gaiitsgory: unramified Hecke operators can slide around,
act centrally on “pinned” ramified Hecke operators



⁸Gukov-Witten

From Physics to Number Theory

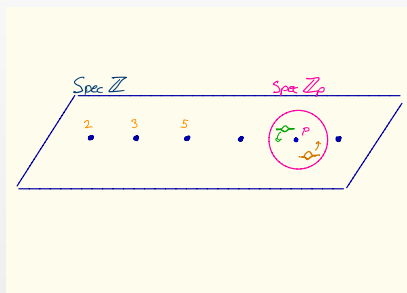
Scholze magically ported factorization to number theory:
geometric source of commutativity of classical Hecke operators



\leadsto Fargues' conjecture: a radically new geometric understanding of the classical (local) Langlands program. cf. Fargues' talk Friday

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Where next?

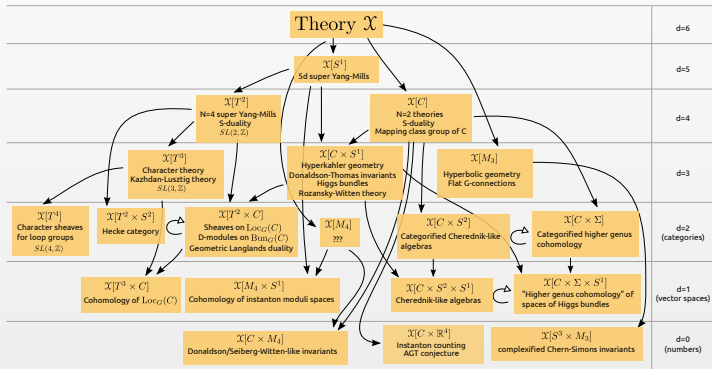
Onward to six dimensions:

Topic of great current interest -
a mysterious six-dimensional QFT, [Theory X](#)

Encodes electric-magnetic duality,
Seiberg-Witten theory, much else in purely geometric terms.

Suggests countless conjectures and inspirations for future work.

The End?



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Schrödinger with Cat:

<http://www.shardcore.org/shardpress/2011/08/05/schrodinger-and-the-cat-2011/>

Logarithm:

<http://ocw.mit.edu/courses/mathematics/18-04-complex-variables-with-applications-fall-1999/study-materials/riemann-surfaces-f-z-log-z2-1/>